










**RESEARCH ARTICLE**

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## The Mechanical Nature of the Lithosphere Beneath the Eastern Central Atlantic Hotspots

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**Key Points:**

- The mechanical effects of the Eastern Central Atlantic hotspots are investigated
- The Azores, Cape Verde, Great Meteor, and Madeira hotspots show standard  $T_e$  values
- The Canary hotspot is the notable exception, with a complex pattern of  $T_e$

**Supporting Information:**

Supporting Information may be found in the online version of this article.

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**Abstract** The Eastern Central Atlantic (ECA) region includes the Azores, Canary, Cape Verde, Great Meteor, and Madeira hotspots. These hotspots exhibit a large variety of characteristics and are rooted in the lithosphere ranging in age from newly created at the Mid Atlantic Ridge to Jurassic at the NW Africa Atlantic margin. Therefore, the ECA region represents an excellent scenario to investigate in an integrated way the effects of hotspots on the mechanical structure of oceanic lithosphere. Here, we calculate the effective elastic thickness ( $T_e$ ) of the lithosphere from an analysis of gravity and topography. Azores hotspot is characterized by a  $T_e < 10$  km, whereas the Great Meteor, Cape Verde, and Madeira hotspots have intermediate  $T_e$  (15–30 km) values. In contrast, the Canary hotspot is characterized by a much higher  $T_e$  (>50 km), forming the largest and most prominent mechanical feature in the ECA. All the hotspots except Canary show standard elastic thickness values when compared to average values for the same age lithosphere and to other oceanic areas in the world. The high strength of the Canary hotspot may be related to the highly depleted mantle composition in the area. The comparison between the elastic thickness distribution and the upper mantle seismic velocity structure shows no correlation between the  $T_e$  estimated at the ECA hotspots (with the exception of Azores) and the presence of low shear-wave velocity anomalies in the underlying mantle. This lack of correlation suggests a negligible effect of upper mantle temperature anomalies on the flexure of the ECA region.

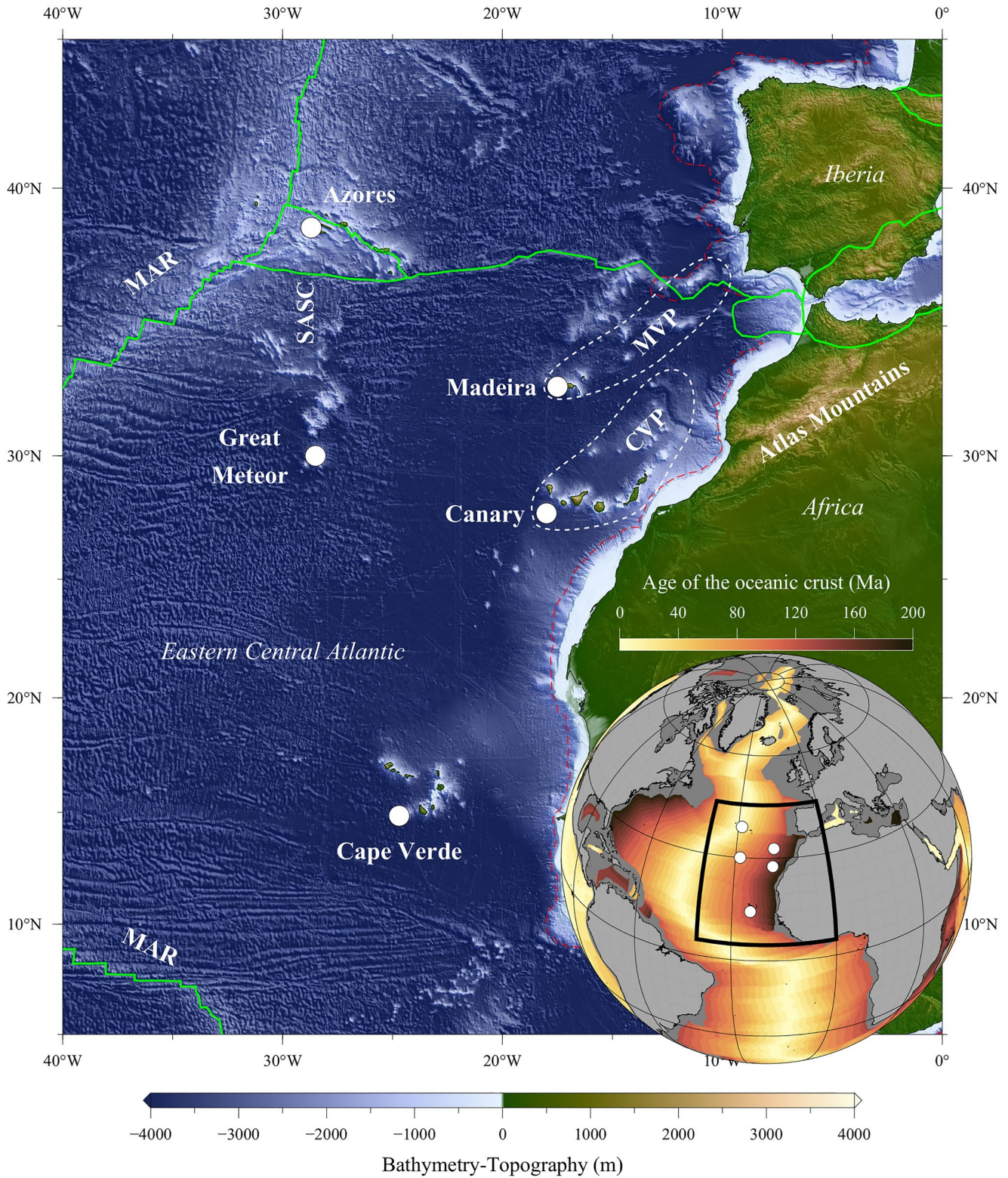
**Plain Language Summary** The upwelling of hot mantle material can modify drastically the thermal and mechanical properties of the rigid outer layer of the Earth (the lithosphere). In this study, we investigate the mechanical effects of the Azores, Canary, Cape Verde, Great Meteor, and Madeira hotspots on the oceanic lithosphere in the Eastern Central Atlantic. To do so, we use gravity and elevation data to analyze the spatial distribution of the effective elastic thickness of the lithosphere in this region, which is a proxy for the lithospheric strength. All of these hotspots except Canary show standard elastic thickness values when compared to the average thickness of the same age lithosphere. In contrast, the lithosphere in the Canary region is much stronger, and this high strength may be related to the nature of the underlying mantle. Finally, the presumed presence of hot mantle beneath these hotspots does not produce any weakening of the overlying lithosphere.

### 1. Introduction

The Eastern Central Atlantic (ECA) is characterized by the presence of a number of volcanic archipelagos, seamounts and plateaux, which are rooted in a crust ranging in age from 0 Ma at the Mid-Atlantic Ridge (MAR) to 180 Ma at the north-western African margin. The ECA includes five hotspots, namely, the Azores, Madeira, the Canary and Cape Verde archipelagos (often collectively termed as Macaronesia region), and the Great Meteor Seamount hotspot (Figure 1) that exhibit contrasting morphologic, geologic, and geophysical characteristics.

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**Figure 1.** Topography and bathymetry of the Eastern Central Atlantic region (GEBCO\_2022 Grid; GEBCO Compilation Group, 2022), showing seafloor age data (Seton et al., 2020) and plate and ocean–continent boundaries (green and red lines, respectively) from Hasterok et al. (2022) for context. The Madeira and Canary volcanic provinces are also shown bounded by white dashed lines (Geldmacher et al., 2005). Topography shaded relief superimposed. Abbreviations: CVP, Canary Volcanic Province; MAR, Mid-Atlantic Ridge; MVP, Madeira Volcanic Province; SASC, Southern Azores Seamount Chain; WAC, West African Craton. The five hotspots (white circles) are also labeled: Azores, Canary, Cape Verde, Great Meteor, and Madeira. The hotspot names and locations are from Jackson et al. (2021).

Therefore, ECA represents an excellent region for investigation of the mechanical effects of hotspots on the lithosphere and the mechanisms supporting intraplate oceanic chains in general.

A useful proxy used to characterize the long-term lithospheric strength is the effective elastic thickness ( $T_e$ ).  $T_e$  corresponds to the thickness of an idealized purely elastic plate of homogeneous properties that would bend similarly to the actual lithosphere under the same applied loads (Watts, 2001), and is related to the integrated mechanical strength of the lithosphere (Burov & Diament, 1995; Watts & Burov, 2003). Previous studies have estimated  $T_e$  by means of different methods in different geological backgrounds. Such methods include direct measurements of flexure from seismic refraction data (e.g., Watts et al., 1985), stratigraphic data (e.g., Collier & Watts, 2001; Grotzinger & Royden, 1990), spectral analysis of the relationship between the topography/bathymetry and gravity (Bouguer or free air) or geoid anomaly (e.g., Audet, 2014; Kirby, 2014; McKenzie & Bowin, 1976) and thermo-rheological modeling of the lithospheric structure (e.g., Cloetingh et al., 2005; Tesauro et al., 2015). The variety of methods has led to different estimates of  $T_e$  even for the same areas, which makes it difficult to compare  $T_e$  results from different studies. This is the case of the five hotspots included in the study area. Furthermore, since many oceanic  $T_e$  studies have been restricted to particular geological features, the mechanical structure of most of the oceanic lithosphere in the study area remains unmapped, and therefore, the possible mechanical effects of hotspots are unknown.

In this work, we present the spatial variations of  $T_e$  in the ECA region from an analysis of the coherence between topography and Bouguer gravity anomaly using wavelet transforms. In order to gain a better insight into the possible origin of the hotspots in the area, we carried out a joint interpretation of the  $T_e$  distribution and the upper mantle seismic structure based on surface-wave tomography. In the following sections, we first describe key observations for the five hotspots included in the study area (Section 2). We then describe the methodology, data and parameters employed for estimating  $T_e$  (Section 3) and finally, we present the spatial variations of  $T_e$  (Section 4) and discuss the implications of our results (Section 5).

## 2. Geodynamic Setting

The Azores Plateau is characterized by a very broad bathymetric anomaly (about 1,000 km wide east-west; Figure 1) containing the triple junction between the North American, Eurasian and Nubian plates. Kinematic reconstructions from 35 Ma (chron 13) propose an age of 20 Ma for the initiation of the formation of the Azores plateau (Gente et al., 2003). The well-developed Azores swell and the presence of low S-wave velocities at 100–200 km depth (e.g., Zhang & Tanimoto, 1993) are consistent with a mantle plume origin for the plateau. However, the lack of any age progression (Feraud et al., 1980) has led to an alternative hypothesis proposing a shallower origin for the plateau related to plate boundary forces (Bonatti, 1990; Sartori et al., 1994).

The Great Meteor Seamount is rooted in a crust of 85 Ma age (Verhoef, 1984), and belongs to the Southern Azores Seamount Chain (SASC) located south of the Azores Plateau (Figure 1). The SASC is characterized by narrow isotopic variability that falls within the Azores isotopic field (Ribeiro et al., 2017). New K-Ar age determinations show nearly coeval volcanism along the chain, with ages of ~31–34 Ma (Ribeiro et al., 2017). A genetic link between the Azores Archipelago and the SASC has been proposed by Ribeiro et al. (2017) on the basis of geophysical observations (Gente et al., 2003) and recent geochemical and isotopic (Sr-Nd-Pb-Hf) data. These authors suggest that the geochemical similarity between the SASC and the Azores mantle source resulted from the long-lived (~85 Ma) Azores Plume.

Crude progression of increasing volcanic ages from southwest to northeast has long been recognized both in the Madeira and the Canary Volcanic Provinces (Abdel-Monen et al., 1972; Geldmacher et al., 2000, 2001; McDougall & Schmincke, 1976). The Madeira Province forms a ~700 km by ~200 km NE-SW-wide chain of volcanoes, with increasing ages from Madeira (0–5 Ma) to the Ormonde Seamount (65–67 Ma). The Canary Province is twice as wide E-W (~700 by ~400 km) and volcanic ages similarly range from 1.1 Ma at El Hierro (Guillou et al., 1996) to 68 Ma in the Lars Seamount (Geldmacher et al., 2001). This trend of increasing age to the northeast is further confirmed by the estimated oldest ages of shield stage volcanism in the seamounts in both provinces (Geldmacher et al., 2005). Although age progression is consistent with the conventional model of a hotspot created by a deep-seated and static mantle plume as first proposed for the Canary Province by Holik et al. (1991), this model neither explains the complex spatial and age distribution of volcanoes, nor the exceptionally long time interval (up to 40 Myr) of active volcanism at single volcanic centers. The cluster of seamounts

southwest of the Canary Islands have been recently dated (van den Bogaard, 2013) and found to be much older than the islands, with ages spatially scattered and ranging between 91 and 142 Ma. Given that these ages do not fit with the trend of increasing age to the northeast, the origin of these seamounts is commonly attributed to fracture-controlled volcanism, rather than to an upwelling mantle plume (van den Bogaard, 2013; Zaczek et al., 2015). The crudeness of the age progression of Madeira and Canary volcanic provinces is in contrast with the systematic and linear Hawaiian age progression derived from the ages of Hawaiian-Emperor seamounts (e.g., Duncan & Keller, 2004).

The Cape Verde Islands have been emplaced on the Early Cretaceous oceanic lithosphere of anomalously shallow bathymetry. The height and width of this swell are 2.2 and 1,200 km, respectively (Ali et al., 2003; Monnereau & Cazenave, 1990). There is no clear age progression of the islands, with estimated ages of about 8 Ma for the western islands and 10–20 Ma for the easternmost islands (Ali et al., 2003 and references therein). Despite the absence of a hotspot track, the hypothesis of an upwelling mantle plume is consistent with the elevated height of the bathymetric swell and with the velocity anomalies of several global and local seismic tomography studies (e.g., Boschi et al., 2007; Celli, Lebedev, Schaeffer, & Gaina, 2020; Celli, Lebedev, Schaeffer, Ravenna, & Gaina, 2020; Liu & Zhao, 2021; Montelli et al., 2006; Schaeffer & Lebedev, 2013; Zhao, 2007; see Section 5).

### 3. Method and Data

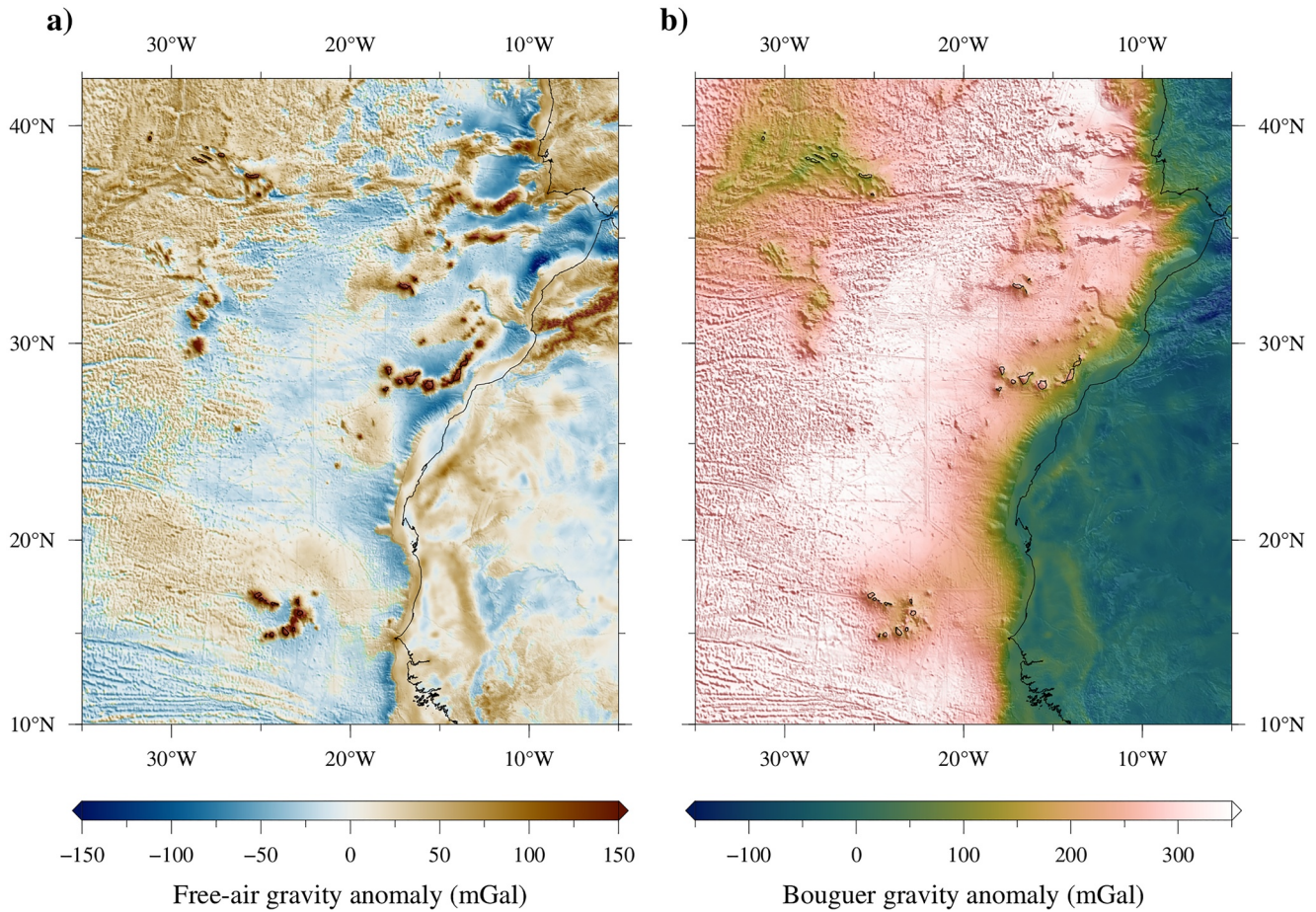
A useful proxy for the long-term lithospheric strength is given by the effective elastic thickness,  $T_e$ , which is related to the flexural rigidity,  $D$ , of a thin elastic plate by

$$D = ET_e^3 / 12(1 - \nu^2) \quad (1)$$

where  $E$  is Young's modulus and  $\nu$  is Poisson's ratio (e.g., Watts, 2001). To estimate  $T_e$  we calculate the coherence function relating the topography and Bouguer anomaly, commonly known as Bouguer coherence (e.g., Audet, 2014; Kirby, 2014), modeled with a simple thin elastic plate subject to both surface and subsurface loads. Specifically, we estimate the gravity and topography spectra, and thence their coherence, using the fan-wavelet transform (Kirby & Swain, 2004, 2011), which provides an observed coherence estimate at every point on the data grid. The wavelet transforms of the grids were computed using space domain convolution at the largest wavelet scales (longest wavelengths) and the Fourier transform at all others (Kirby & Swain, 2013). We use Morlet wavelet central wavenumber values of  $|\mathbf{k}_0| = 2.668$  and 5.336, where the smallest value prioritizes the spatial resolution at the expense of wavenumber resolution, and vice versa for the largest value (Kirby & Swain, 2011).

The observed wavelet squared real coherency (SRC) is inverted for  $T_e$  using fan-wavelet adaptation of Forsyth's (1985) method (Kirby & Swain, 2008, 2009; Swain & Kirby, 2006). We follow Kirby and Swain (2009) and invert the SRC since it is less sensitive than the coherence to correlations between the initial loads on the plate and to the presence of topographically unexpressed internal loading (commonly known as “gravitational noise”), both of which can cause incorrect recovery of  $T_e$  (e.g., Audet & Bürgmann, 2011; Kirby, 2014; Kirby & Swain, 2009; McKenzie, 2003; McKenzie & Fairhead, 1997). In general, subsurface (internal) loads include mafic intrusions, accreted lower crustal material, thermal anomalies, and compositional variations, which cause lateral variations of density at depth, while surface loading is caused by topography and large-scale variations in surface density (e.g., mountains and sedimentary basins). Forsyth's (1985) method also yields the initial load ratio ( $f$ ) between initial internal and surface loads, which we express here as an internal load fraction,  $F$ , where  $F = f/(1 + f)$  (McKenzie, 2003). Pure surface loading gives  $F = 0$ , purely internal loading gives  $F = 1$ , while equal surface and internal loading gives  $F = 0.5$ . We use Brent's method (Press et al., 1992) in the inversion to find the best-fitting predicted SRC, with the difference between observed and predicted SRCs weighted by the inverse of the equivalent Fourier wavenumber (Kirby & Swain, 2006). This inverse-wavenumber method generally gives smoother  $T_e$  maps than with computed errors, as it down-weights the noisy SRCs that sometimes occur at high wavenumbers. Errors on  $T_e$  are given as 95% confidence limits on the best-fitting value. For an extensive description of the method, see the Supplementary Material.

The bathymetry/topography data (Figure 1) used in our analysis are obtained from the GEBCO\_2022 Grid, a continuous, global terrain model for ocean and land with a spatial resolution of 15 arc seconds (GEBCO Compilation Group, 2022). Although GEBCO 2022 uses SRTM15+ as a base, which contains altimetry-derived bathymetry, it is extensively supplemented with ship echo soundings, making it a much more quasi-independent dataset than bathymetry grids derived exclusively from altimetry estimates of the geoid. Therefore, the GEBCO



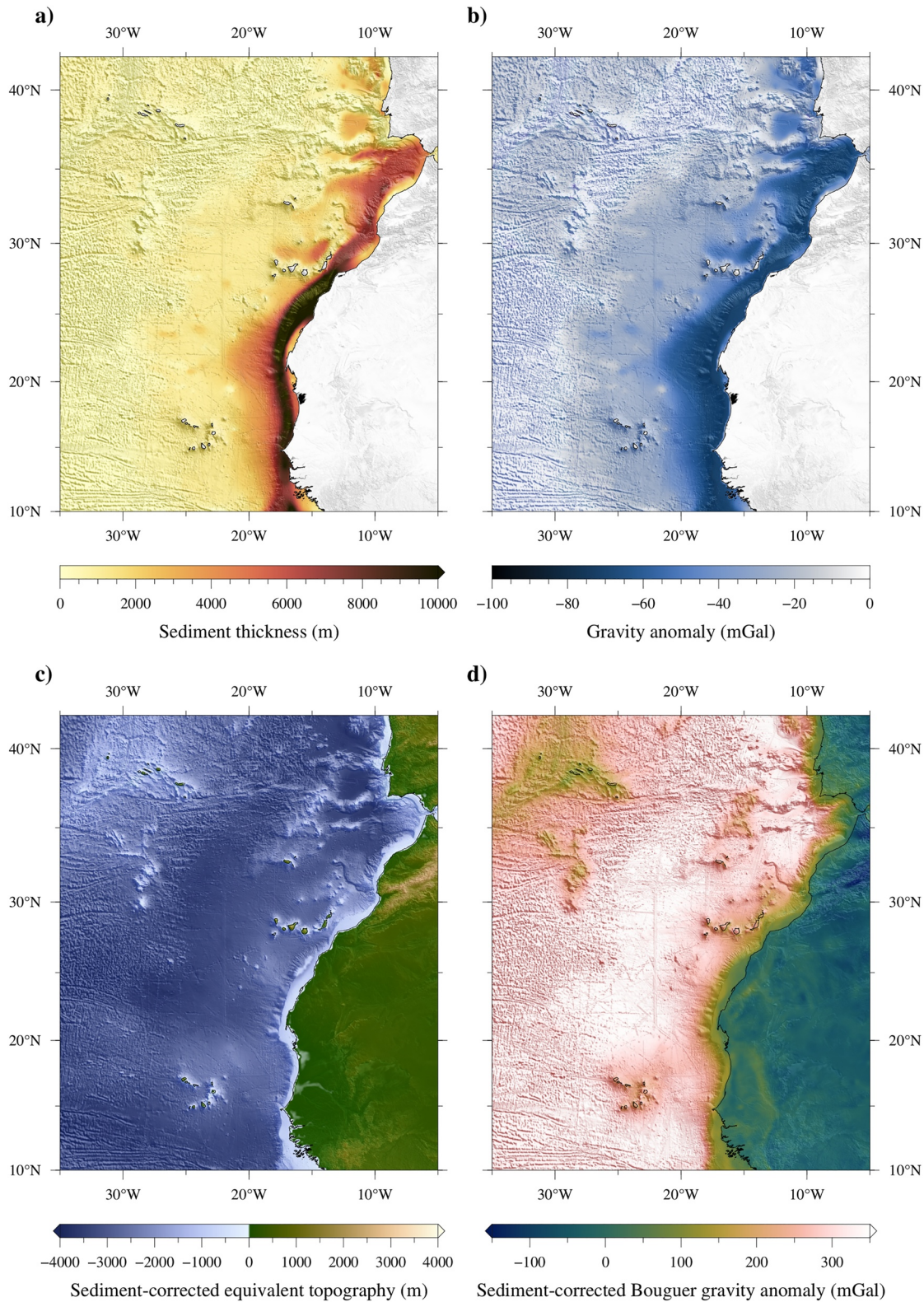
**Figure 2.** (a) Free-air and (b) Bouguer gravity anomalies used for the analysis (EGM2008 gravity model; Pavlis et al., 2008). Topography shaded relief superimposed.

bathymetry is preferred in  $T_e$  estimation over the oceans because it maximizes the independence of the gravity and bathymetry data sets (e.g., Kalnins & Watts, 2009; Lu et al., 2021). Our area contains both continental and oceanic lithosphere, with the latter being subject to an additional water load. To treat mixed land and marine environments, we adopt the approach of Stark et al. (2003) and Kirby and Swain (2008). This approach scales ocean bathymetry ( $h$ ) to equivalent topography prior to Fourier transformation:

$$h_{\text{eq}} = h(\rho_w - \rho_w) / \rho_{\text{uc}} \quad (2)$$

where  $\rho_w$  and  $\rho_{\text{uc}}$  are the seawater density ( $1,030 \text{ kg m}^{-3}$ ) and upper crust density ( $2,670 \text{ kg m}^{-3}$ ), respectively; that is, we scale the load of the ocean water to an equivalent topographic load. This allows the loading equations for a land environment to be used for the whole area, rather than performing two separate analyses and inversions on land and ocean areas, avoiding  $T_e$  discontinuities at the coastlines (e.g., Jiménez-Díaz et al., 2014; Pérez-Gussinyé et al., 2004; Ratheesh-Kumar & Xiao, 2018; Shi et al., 2017). For the purposes of comparison,  $T_e$  calculated using the regular bathymetry/topography is also presented in the Supplementary Material (Figures S1–S3 in Supporting Information S1). The Free-air and Bouguer gravity anomaly data over both land and ocean are obtained from the EGM2008 gravity model (Pavlis et al., 2008), with a spatial resolution of 2.5 arc minutes (Figure 2).

Sediments may play an important role in  $T_e$  estimation using spectral methods due to the significant density contrast with the underlying basement, which is not considered a priori (e.g., Chen et al., 2015; Ji et al., 2020; Kaban et al., 2018; Kirby, 2022; Ratheesh-Kumar et al., 2015; Shi et al., 2017; Yu et al., 2022). Large amount of sediments cause a flattening of the bathymetry and a prominent negative gravity anomaly, modifying the relationship between surface topography and gravity anomalies. Therefore, this combined effect of sediments can bias the  $T_e$  estimates based on the correlation between topography and gravity. Figure 3a shows the seafloor sediment thickness in the study area. The sediment thickness data are obtained from the GlobSed global model



**Figure 3.** (a) Sediment thickness (GlobSed global model; Straume et al., 2019). (b) Gravity effect of sediments. (c) Sediment-corrected equivalent topography (derived from the GEBCO\_2022 Grid; see the main text) and (d) sediment-corrected Bouguer gravity anomaly (derived from the EGM2008 gravity model) used for the analysis. Topography shaded relief superimposed.

(Straume et al., 2019), an updated total sediment thickness grid for the world's oceans and marginal seas with a spatial resolution of 5 arc minutes. Large amount of sediments are mainly located on the Gulf of Cadiz basin and offshore areas of western Iberia, and especially along the western continental margin of Africa where, in many cases, the sediment thickness exceeds 10 km. This can produce significant negative gravity anomalies reaching values of up to  $-80$  mGal, as can be seen in Figure 3b.

Therefore, we examine the spatial variations of  $T_e$  in the study area considering sediments. We do so by using sediment-corrected bathymetry and gravity data sets. The gravity effect of sediments (Figure 3b) is calculated using the method of Rao and Babu (1991), and the sediment density model refers to the density-depth curves of offshore basins based on Mooney and Kaban (2010). The final Bouguer gravity anomaly is then calculated by subtracting the gravity effect of sediments from the complete Bouguer gravity anomaly. In addition, following the method of converting ocean bathymetry to equivalent topography (see above), the sediment thickness is also corrected, obtaining the final equivalent topography:

$$h_{eq} = h(\rho_w - \rho_{uc})/\rho_w + s(\rho_s - \rho_{uc})/\rho_w \quad (3)$$

where  $\rho_w$  and  $\rho_{uc}$  are the seawater density ( $1,030 \text{ kg m}^{-3}$ ) and upper crust density ( $2,670 \text{ kg m}^{-3}$ ), and  $\rho_s$  and  $s$  are the sediment density and thickness, respectively; also scaled according to the sediment's density-depth curves (Mooney & Kaban, 2010). The sediment-corrected equivalent topography and Bouguer gravity anomaly are shown in Figure 3. To be clear, note that the flexural component of the sediment effect on bathymetry is not corrected for. The flexural effect of sediments depends on  $T_e$ , which is unknown a priori. We treat them as just another layer in the crust, and replace them with rock of mean crustal density (we actually account for a depth-dependent density; see above). It is, in essence just a Bouguer correction and an equivalent topography correction, but instead of replacing seawater with rock (or ice with rock), we replace sediments with rock, and work out the (a) change in gravity, and (b) the change in bathymetry. This correction considerably suppresses the effect of unexpressed subsurface loads and substantially reduces  $T_e$  estimates in areas where the topography variations are negligible and density variations within the sedimentary cover dominate (e.g., Kaban et al., 2018; Shi et al., 2017; Yu et al., 2022).

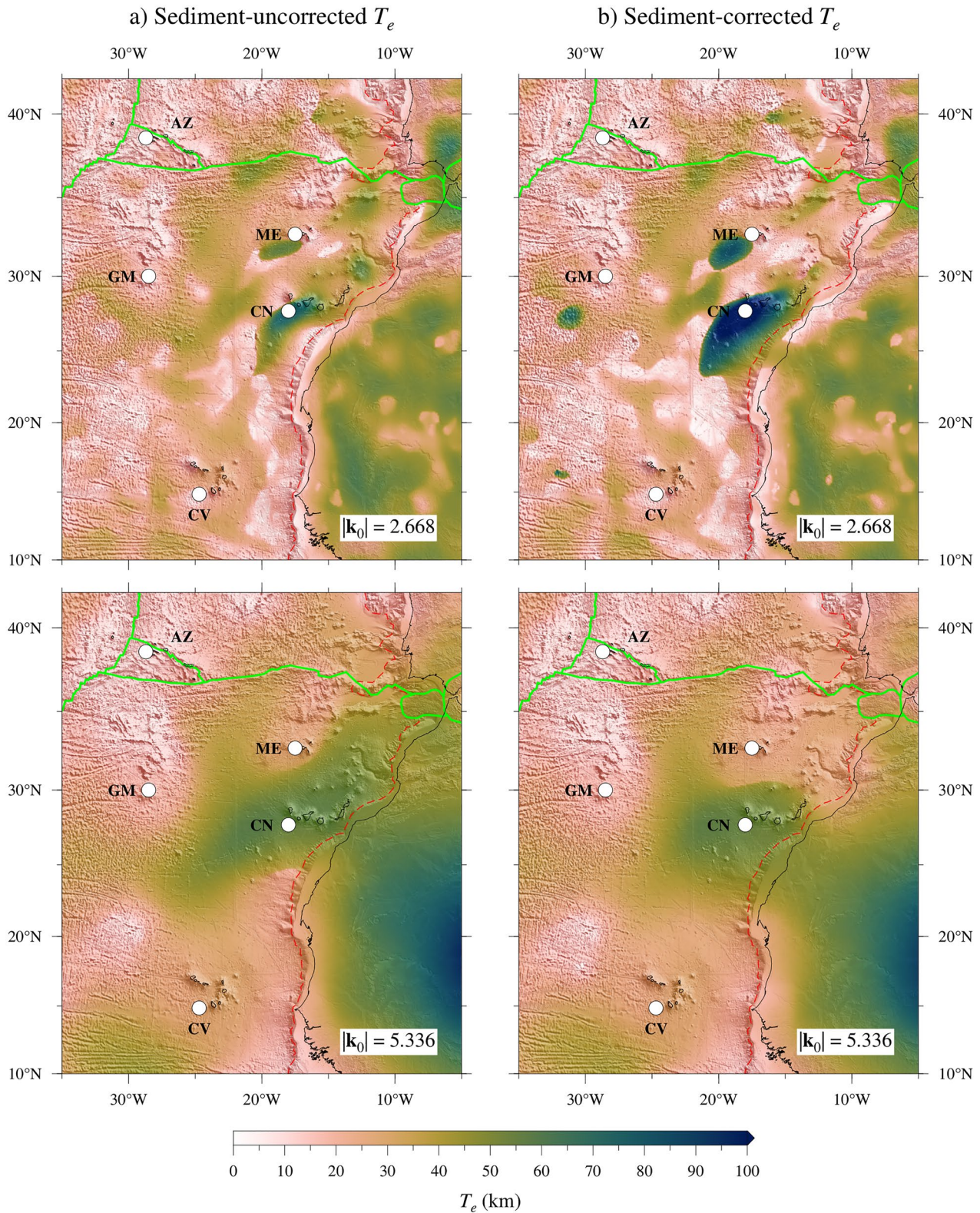
Forsyth (1985)'s method (often called "load deconvolution" as it seeks to recreate the initial loads from the final, post-flexure gravity and topography) requires estimates of crustal structure. Here, we assume a single-layered crust with internal loading occurring at the Moho, as did Forsyth (1985), and use the Moho depths and crustal densities provided by the global crustal model CRUST1.0 (Laske et al., 2013). Other constants used in the SRC inversion are the following: Newtonian gravitational constant,  $G = 6.67259 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$ ; Young's modulus,  $E = 100 \text{ GPa}$ ; Poisson's ratio,  $\nu = 0.25$ ; the gravity acceleration,  $g = 9.79 \text{ m s}^{-2}$ ; and mantle density,  $\rho_m = 3,300 \text{ kg m}^{-3}$ .

We project all data sets to a Cartesian coordinate system using the Mercator projection in order to mitigate errors arising from the planar treatment of curvilinear coordinates. The study area has dimensions of  $4,040 \text{ km}$  (east-west)  $\times$   $4,570 \text{ km}$  (north-south), and a grid spacing of  $10 \text{ km}$  in both directions. The gravity/topography data are mirrored about their edges prior to Fourier transformation with the purpose of reducing spectral leakage, which in combination with the wavelet transform barely bias the results (see Kirby & Swain, 2008).

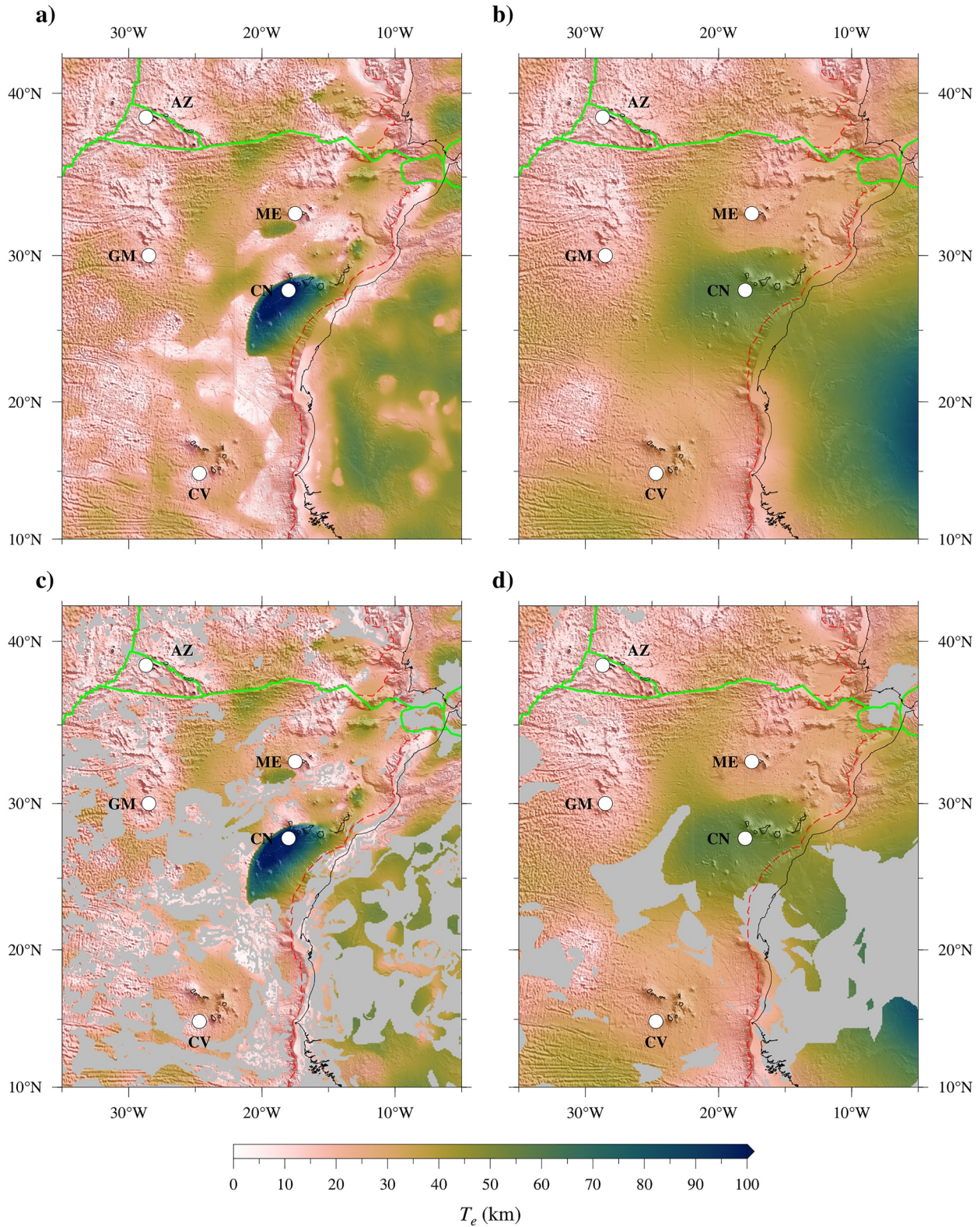
#### 4. Results

The effective elastic thickness results for the two  $|\mathbf{k}_0|$  values are shown in Figures 4–6. Low values of  $|\mathbf{k}_0|$  provide better spatial resolution, while large values give a good resolution in the wave-number domain and a reasonably good resolution in the space domain. In this way, the results with different  $|\mathbf{k}_0|$  can be interpreted as the result of applying different band-pass filters to  $T_e$ , bearing in mind that absolute  $T_e$  values are more reliable for higher  $|\mathbf{k}_0|$  wavelets, while relative  $T_e$  differences are more faithfully represented by lower  $|\mathbf{k}_0|$  wavelets (Kirby & Swain, 2011; Swain & Kirby, 2021). All of this can be seen in Figure 4, where the obtained  $T_e$  map with the  $|\mathbf{k}_0| = 2.668$  shows a more detailed  $T_e$  structure, whereas  $T_e$  results from  $|\mathbf{k}_0| = 5.336$  show a spatially smoother  $T_e$  pattern within the region. Thus, combining different  $|\mathbf{k}_0|$  offers a comprehensive overview of  $T_e$  variations for the study area at different spatial scales.

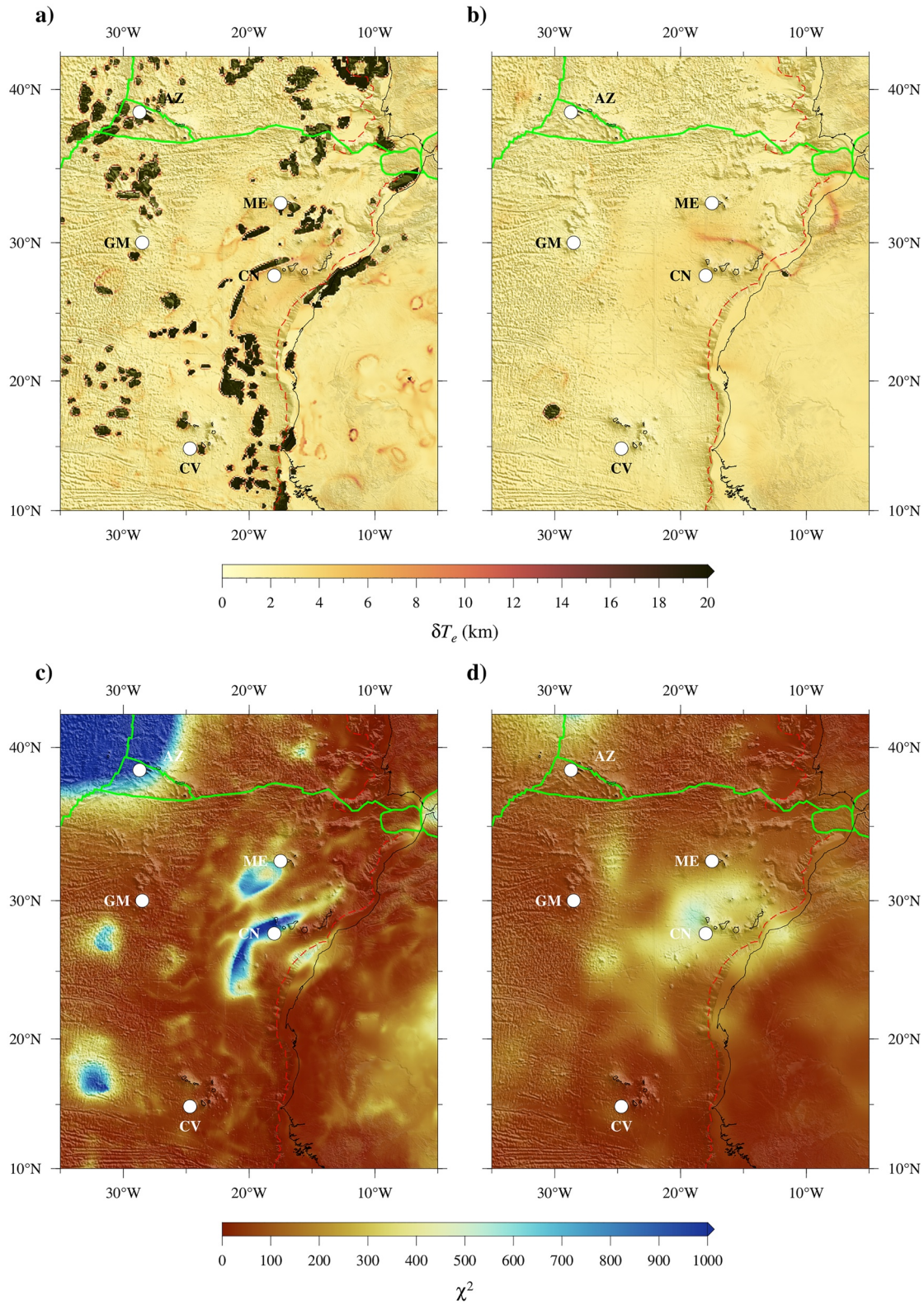
For the purposes of comparison, Figures 4a and 4b show the sediment-uncorrected and the sediment-corrected  $T_e$  results estimated with  $|\mathbf{k}_0| = 2.668$  and  $5.336$ , respectively. In general, the  $T_e$  results using the lowest  $|\mathbf{k}_0|$  ( $2.668$ ) wavelet remain similar, though the Canary Volcanic Province is characterized by a comparatively much



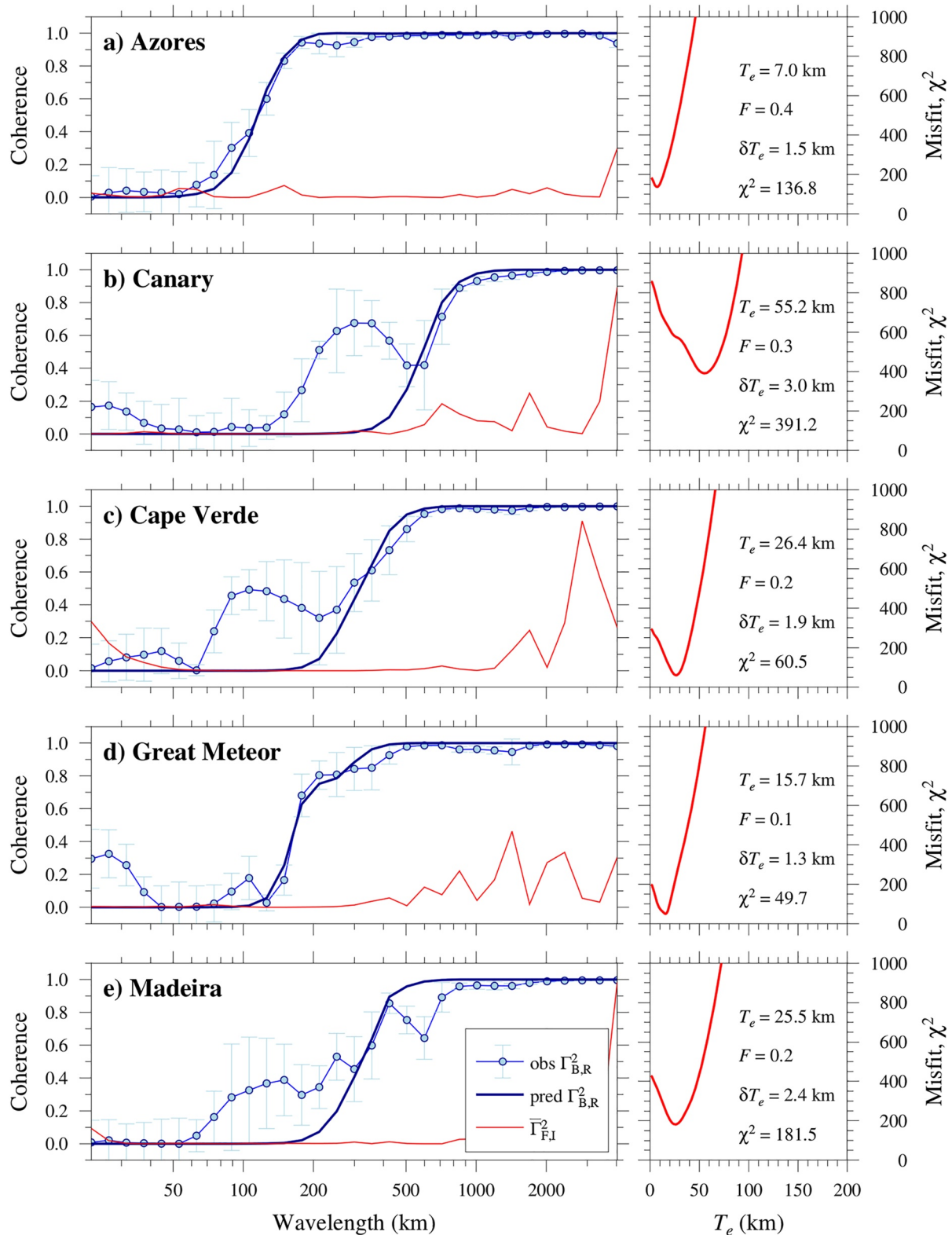
**Figure 4.** Effective elastic thickness,  $T_e$ , from the inversion of the squared-real Bouguer coherency using (a) the sediment-uncorrected and (b) the sediment-corrected equivalent topography and Bouguer gravity anomaly, with  $|k_0| =$  (a) 2.668 and (b) 5.336, respectively. Abbreviations: AZ, Azores; CN, Canary; CV, Cape Verde; GM, Great Meteor; ME, Madeira. The hotspot names, abbreviations, and locations are from Jackson et al. (2021).



**Figure 5.** (a and b) Effective elastic thickness,  $T_e$ , from the inversion of the squared-real Bouguer coherency using the sediment-corrected equivalent topography and Bouguer gravity anomaly, for  $|k_0| =$  (a) 2.668 and (b) 5.336, respectively. (c and d) Effective elastic thickness,  $T_e$ , corresponding to the results in (a and b). Gray-shaded areas correspond to regions where  $T_e$  estimation might be biased by "gravitational noise." White circles indicate the locations where squared real coherency profiles ( $|k_0| = 5.336$ ) are plotted in Figure 7. Abbreviations are the same as in Figure 4.



**Figure 6.** (a and b)  $T_e$  error and (c and d) chi-square misfit, corresponding to the  $T_e$  results in Figure 5, for  $|k_0| =$  (a) 2.668 and (b) 5.336, respectively. The color scale in (a) saturates at 20 km. Abbreviations are the same as in Figure 4.



**Figure 7.** (a–e) Example of curve fitting and inversion results using the Bouguer squared real coherency (SRC) ( $\Gamma_{B,R}^2$ ) for the five hotspots with  $|\mathbf{k}_q| = 5.336$ . The autospectra and cross-spectra were averaged over spatial dimensions of  $150 \times 150$  km and the observed coherency formed from these. Left panels: The observed Bouguer SRC (blue circles with their error bars and blue lines); the best-fitting predicted Bouguer SRC (dark blue lines); and the normalized free-air squared imaginary coherency,  $\Gamma_{F,I}^2$  (thin red lines). Right panels: misfit curves (thick red lines) used to estimate  $T_e$  using the Bouguer SRC; and the values of the best fitting  $T_e$  and  $F$ , as well as the value of the  $T_e$  error and chi-square misfit at all locations.

higher  $T_e$  than previously observed in Figure 4a. In contrast, we note that the  $T_e$  results for  $|\mathbf{k}_0| = 5.336$  show some important differences when considering the sediments. The most noticeable differences occur within the Gulf of Cadiz and the Madeira Volcanic Province, and along the northwestern continental margin of Africa, where the sediment-corrected  $T_e$  values are significantly smaller than the  $T_e$  results in Figure 4a, suggesting that the thick and less dense sediments have a large effect on the  $T_e$  results in these areas. Importantly, the presence of sediments does not significantly affect our  $T_e$  estimates for the hotspots in terms of absolute values (see for comparison Figure 7 and Figure S6 in Supporting Information S1). In the following section, we analyze our sediment-corrected  $T_e$  results in detail. Sediment-uncorrected results are also presented in Figures S4–S6 in Supporting Information S1.

Figure 5 shows the sediment-corrected  $T_e$  results with  $|\mathbf{k}_0| =$  (a) 2.668 and (b) 5.336, respectively. Azores is characterized by a  $T_e < 10$  km for both  $|\mathbf{k}_0|$  values. The Great Meteor and Cape Verde show a slightly higher  $T_e$  value of 15–30 km, both being bounded by higher  $T_e$  values (especially noticeable with the lowest  $|\mathbf{k}_0|$  value; see Figure 5a). The Madeira and Canary volcanic provinces are characterized by a complex pattern of  $T_e$ . The Madeira Islands exhibit a  $T_e$  of 15–30 km for both  $|\mathbf{k}_0|$  values. However, for  $|\mathbf{k}_0| = 2.668$ , this  $T_e$  pattern over the Madeira Islands shows a moderate increase along the northeastern Madeira Volcanic Province, and an abrupt increase in  $T_e$  along the southwest of the Madeira Islands (Figure 5a). Finally, the Canary Islands are characterized by  $T_e$  values exceeding 50 km. This  $T_e$  pattern also extends to the west and southwest of the Canary Islands (with a steep gradient with the lowest  $|\mathbf{k}_0|$  value; see Figure 5a), forming the largest, and highest in terms of  $T_e$ , feature in the Eastern Central Atlantic. We note that  $T_e$  results for  $|\mathbf{k}_0| = 5.336$  appear to be more influenced by the high values recovered over Africa, where the high  $T_e$  of the West African Craton has been extended westwards (due to the comparatively lower spatial resolution of such wavelets).

We analyze the bias in  $T_e$  estimation related to the effect of gravitational noise following the approach of Kirby and Swain (2009), based on the normalized free-air squared imaginary coherency (SIC). Figures 5c and 5d show the effective elastic thickness,  $T_e$ , from the inversion of the squared-real Bouguer coherency (“Bouguer SRC”) with  $|\mathbf{k}_0| =$  (a) 2.668 and (b) 5.336, respectively. Gray-shaded areas correspond to regions where  $T_e$  estimation might be biased by noise based on our analysis (i.e., regions where the maximum value of the normalized free-air SIC exceeds 0.5). Note that while large noise indicates a potential failure of the load deconvolution method, it does not imply that the lithosphere is weak (Audet & Bürgmann, 2011). Kirby and Swain (2011) found that the size of the affected area increases as  $|\mathbf{k}_0|$  decreases (see Figures 5c and 5d), and also that the lower wavenumber resolution of low- $|\mathbf{k}_0|$  wavelet results in the noise spectrum being smeared over a larger bandwidth than with high- $|\mathbf{k}_0|$  wavelets. In this regard, we note that noise does exist in the study area. However, the five hotspots are either slightly affected (Azores and Great Meteor; see Figure 5) or even unaffected, especially when using the highest  $|\mathbf{k}_0|$  value.

Figure 6 shows  $T_e$  error and the spatial variation of the chi-square statistic minimized in the inversion of the Bouguer SRC ( $|\mathbf{k}_0| = 2.668$  and 5.336), corresponding to the  $T_e$  results in Figure 5. We limit the color scale to a maximum  $T_e$  error of 20 km. Uncertainty for  $T_e$  estimates is low, although higher errors (with a maximum  $T_e$  error >100 km) are found in some areas surrounding the five hotspots in the lowest  $|\mathbf{k}_0|$  results (Figure 6a). We note that the chi-square misfit can be quite high in some regions, especially when using the lowest  $|\mathbf{k}_0|$  (Figure 6c). Note that Figures 6c and 6d provide a measure of how well the observed and best-fitting predicted SRC fit. Therefore, Figures 6c and 6d contain qualitative information about  $T_e$  errors and show regions where the model may be inappropriate and/or the data errors have been underestimated. We find the highest values in the northwest of the Azores Plateau. Other areas with very high chi-square misfit values coincide with the higher  $T_e$  along the southwest of the Madeira Islands (see Figure 5a), and a narrower band along the southwest of the Canary Islands, which likely reflects an overestimation in these areas.

Finally, as an additional test of the robustness of the inversion and  $T_e$  estimates, Figure 7 shows the results of inverting wavelet SRC profiles for the five hotspots with  $|\mathbf{k}_0| = 5.336$ , hence prioritizing absolute  $T_e$  values. Rather than computing the coherency at each of the  $10 \times 10$  km grid nodes, we spatially averaged the auto-spectra and cross-spectra over  $15 \times 15$  grid nodes, giving a  $150 \times 150$  km spatially averaged coherency. The thin red lines show the normalized squared imaginary component of the free-air coherency. Large values at wavelengths just less than the transition wavelength (from low to high Bouguer SRC) indicate a  $T_e$  estimate likely biased by gravitational noise. High values of the normalized free-air SIC occur at wavelengths that do not affect the SRC rollover in these regions (Figure 7). Furthermore, the chi-square misfits are all very narrow with well-defined minima, indicating that the recovered  $T_e$  values are robust. However, we note that Canary and Cape Verde present a

complex pattern with two rollovers, suggesting two dominant  $T_e$  values (see also Figures S3 and S6 in Supporting Information S1). For Canary, the longer-wavelength rollover at 600–700 km wavelength provides the estimation. If we consider the shorter-wavelength rollover at  $\sim 200$  km wavelength, a lower  $T_e$  would be plausible. The long, shallow-gradient rollover in Madeira also suggests a wide range of  $T_e$  values, possibly similar to Canary and Cape Verde. These results suggest the support of different loads that may be operating at small- and large-scales simultaneously, or even at different times, in these regions. In this regard, it is worth noting that the volcanic activity is very long lasting, particularly in the Canary volcanic province, and therefore the lithosphere was loaded at different times. For example, small seamounts to the south of El Hierro have volcanic ages as old as 140 Ma (van den Bogaard, 2013), so they were emplaced in much younger lithosphere, which could have responded in a very different way from the Jurassic lithosphere at the time of the formation of the Canary Islands (subaerial volcanic ages 23–1 Ma). We will return to this issue in Section 5.4.

## 5. Discussion

### 5.1. Comparison With Previous $T_e$ Estimates

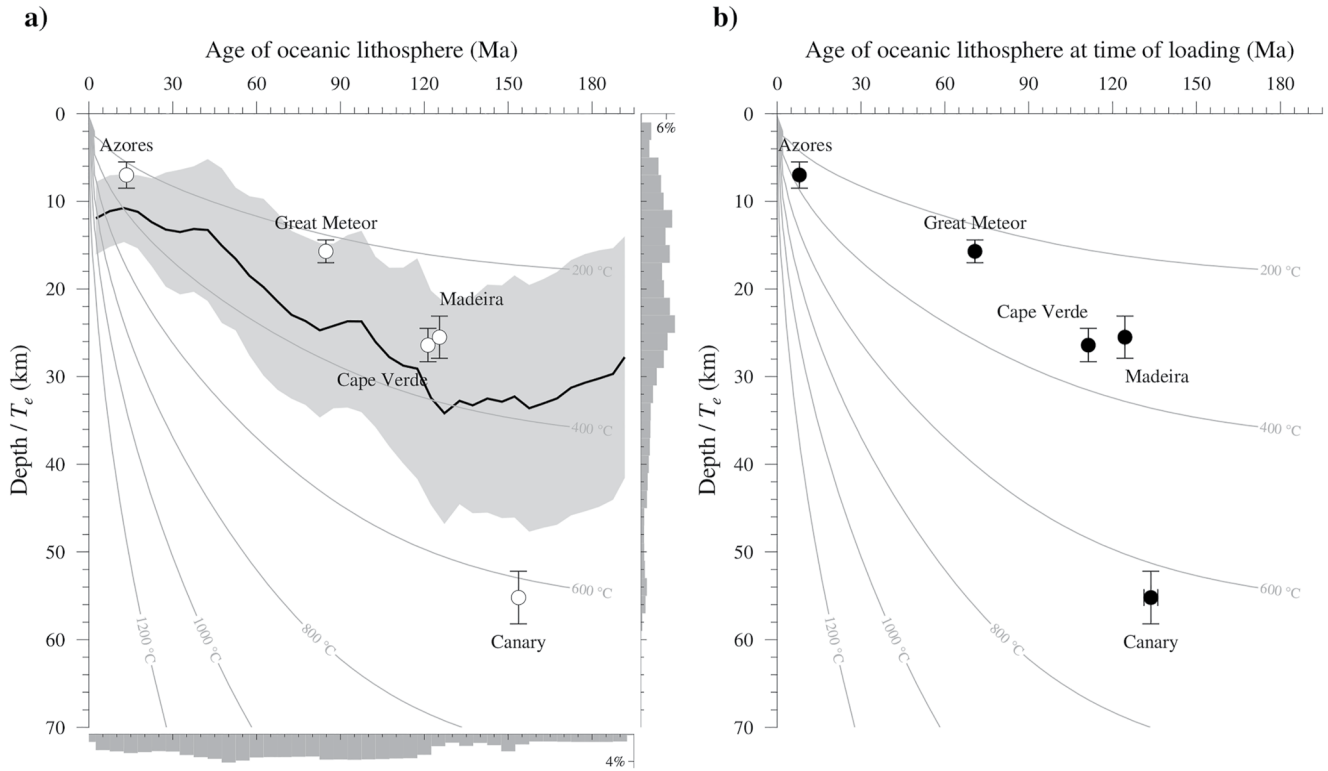
The values obtained for  $T_e$  are, in general, in good agreement with previous estimates for the different analyzed regions, except for the Canary Volcanic Province. The present study predicts a  $T_e$  value for the Azores Archipelago and SASC areas as low as 5–7 km, in good agreement with the  $T_e$  of  $3 \pm 3$  km obtained by Calmant et al. (1990), and of 3–6 km obtained by Luis and Neves (2006) from 3D admittance and coherence study. The obtained  $T_e$  for Great Meteor ( $\sim 16$  km) is consistent with the values of around 20 km obtained by Watts et al. (1975) and Calmant et al. (1990). Similarly, taking into account the uncertainties, our results for Cape Verde (24–28 km) and Madeira (23–28 km) are comparable to previously reported values of, respectively, 12–35 km (Ali et al., 2003; Calmant et al., 1990; Kalnins, 2011; McNutt, 1988; Young & Hill, 1986) and 15–35 km (Calmant et al., 1990; Kalnins, 2011).

In contrast, the  $T_e$  values derived here ( $>50$  km) for Canary are at least 10–35 km higher than the range of 15–40 km obtained from several previous works (Canales & Dañoibeitia, 1998; Collier & Watts, 2001; Dañoibeitia et al., 1994; Kalnins, 2011; Llanes, 2006; Watts, 1994; Watts et al., 1997), but similar to the  $\sim 50$  km estimates by Filmer and McNutt (1989) and consistent with the  $T_e$  maps obtained by Pérez-Gussinyé et al. (2009) and Audet and Bürgmann (2011). The flexural analysis of the static elevation residuals by Fulla et al. (2015) also suggested a mechanically strong lithosphere, in line with the normal thermal lithospheric thicknesses of  $110 \pm 20$  km obtained in their study. This is consistent with the presence of a relatively small amplitude swell (Canales & Dañoibeitia, 1998) and the absence of long-wavelength gravity and geoid anomaly centered in the Canary Islands (Fulla et al., 2015; Watts, 1994). We will return to this issue in Section 5.4.

### 5.2. Relationship Between $T_e$ and the Temperature of the Oceanic Lithosphere

The usual way of systematizing the relation between the mechanical and thermal states of the oceanic lithosphere is to compare  $T_e$  and the age of the lithosphere at the time of loading (i.e., the difference between the age of the lithosphere and the age of the load), with the depth to specific isotherms. Based on that, several works have found a general correlation between  $T_e$  and the depth to the 300–600°C isotherms (e.g., Watts, 2001).

Figure 8a shows the relationship between  $T_e$  and the age of the oceanic lithosphere for the five hotspots studied in this work ( $T_e$  results using  $|k_0| = 5.336$ ; see Figure 7 and Table 1). The isotherm distribution corresponds to the cooling-plate model of Parsons and Sclater (1977). For the purposes of comparison, the thick black line represents the mean  $T_e$  estimates for the same age from the gridded  $T_e$  results using  $|k_0| = 5.336$  (Figure 5), versus the age of the oceanic crust from Seton et al. (2020). There is a systematic increase in  $T_e$  with plate age for lithospheres younger than 130 Ma, consistent with models for the cooling of oceanic lithosphere as it moves away from a mid-ocean ridge and the temperature-dependent ductile creep of lithospheric minerals such as olivine (e.g., Hunter & Watts, 2016). For older lithospheres,  $T_e$  remains nearly constant or even decreases with plate age from 34 to 28 km. Low  $T_e$  values (compared to expected values for very old lithosphere) are found along most passive continental margins. The reason why passive margins tend to be relatively weak remains an open question, with some authors invoking thermal and fluid-related effects in rifting processes (e.g., Audet & Bürgmann, 2011; Pérez-Gussinyé et al., 2009; Ratheesh-Kumar & Xiao, 2018) or accumulated weakening of the passive margins caused by grain damage (Mulyukova & Bercovici, 2018). In turn, Bercovici and Mulyukova (2021) have shown that combined grain damage and mixing enlarges the range of temperatures and depths for the occurrence of this



**Figure 8.** (a)  $T_e$  estimates for the five hotspots versus the age of the oceanic lithosphere. Thick black line represents the mean  $T_e$  estimates ( $\pm$  the standard deviation about the mean; gray envelope) from the gridded  $T_e$  results using  $|k_0| = 5.336$  (Figure 5b), versus the age of the oceanic crust from Seton et al. (2020). The gray curves show the depth to the plate-cooling isotherms of Parsons and Sclater (1977). (b)  $T_e$  estimates for the five hotspots versus the age of the oceanic lithosphere at the time of loading (see Table 1).

weakening. The average  $T_e$  obtained for the ECA region roughly follows the depth of the 350–400°C isotherms, with an envelope of approximately 200–500°C (described by the standard deviation band).

Figure 8b shows the relationship between  $T_e$  and the age of the lithosphere at the time of loading. Although the age of the lithosphere was similar when it was loaded at Canary, Cape Verde and Madeira areas, the  $T_e$  results are very variable, implying a correlation with isotherms depths expanding by roughly a factor two (taking into account uncertainties). The  $T_e$  for Canary roughly correlates with the depth to the isotherm of  $640 \pm 40^\circ\text{C}$ , whereas  $T_e$  for Cape Verde and Madeira correlates with the depth to the isotherm of  $340 \pm 30^\circ\text{C}$  and of  $310 \pm 40^\circ\text{C}$ , respectively. The  $T_e$  for Azores correlates with the depth to the isotherm of  $340 \pm 80^\circ\text{C}$ , and the  $T_e$  obtained for the lithosphere loaded by the Great Meteor seamount correlates with the depth to an isotherm of  $250 \pm 30^\circ\text{C}$ .

**Table 1**

Summary of the Estimates of Elastic Thickness,  $T_e$

	Age of seafloor (Ma) <sup>a</sup>	Age of load (Ma)	Load age error (Myr)	$T_e$ (km) <sup>b</sup>	$T_e$ error (km) <sup>b</sup>
Azores	13.47	5.5 <sup>c</sup>	-	7.0	1.5
Canary	153.68	20 <sup>d</sup>	2.5 <sup>d</sup>	55.2	3.0
Cape Verde	121.31	10 <sup>e</sup>	0 <sup>e</sup>	26.4	1.9
Great Meteor	84.76	14 <sup>e</sup>	0 <sup>e</sup>	15.7	1.3
Madeira	125.41	1 <sup>e</sup>	0 <sup>e</sup>	25.5	2.4

<sup>a</sup>Age of the oceanic crust from Seton et al. (2020). <sup>b</sup> $T_e$  and  $T_e$  error from the results of inverting wavelet squared-coherency profiles ( $|k_0| = 5.336$ ) for the five hotspots (see Figure 7). <sup>c</sup>Based on Feraud et al. (1980). <sup>d</sup>Dañoibeitia et al. (1994). <sup>e</sup>Calmant et al. (1990).

The  $T_e$  value for the Great Meteor can be considered as relatively low when compared with the average value for the lithosphere of the same age in the ECA region, although it falls within the standard deviation band (Figure 8a). The correlating isotherm of  $250 \pm 30^\circ\text{C}$  is slightly colder than the envelope of 300°C–600°C found in the extensive compilation by Watts and Zhong (2000), but it is still within the range of values for a number of seamounts and volcanic islands computed in more recent studies (Kalnins & Watts, 2009; Watts et al., 2006). For example, Watts et al. (2006) compared the predicted bathymetry to shipboard bathymetry measurements and obtained >9,000  $T_e$  estimates from a broad range of submarine features in the Pacific, Indian, and Atlantic oceans included in the Wessel (2001) global seamount database. They found that many  $T_e$  estimates plotted outside the expected 300–600°C isotherm depth range. A re-assessment by Kalnins and Watts (2009) in the area of the Western Pacific Ocean found that bathymetric features formed on the Jurassic seafloor have an envelope of 50–300°C, with a best fit of

180°C. Similarly, the recent study by Lu et al. (2021) found that 80% of their  $T_e$  estimates over the Pacific basin correspond to depths of a broad range of isotherms, with envelopes of  $220 \pm 195^\circ\text{C}$  and  $270 \pm 210^\circ\text{C}$  based on two different cooling plate models. Although the present study further supports that there is no single isotherm that describes all the  $T_e$  data even for the same region (wide standard deviation in Figure 8a), four of the five hotspots analyzed plot within the 250–350°C isotherm, in line with the above-mentioned studies. In contrast, the well-studied Hawaiian chain has a  $T_e$  value of 30–35 km (Watts, 1978; Zhong & Watts, 2013), which correlates with the warmer  $450 \pm 150^\circ\text{C}$  isotherm for a lithospheric age at the time of loading of 70–80 Ma. Overall, the only clearly anomalous  $T_e$  value is the one estimated for the Canary province, with a value  $>50$  km correlating with the  $\sim 650^\circ\text{C}$  isotherm. This correlation with deep geotherms is found only in oceanic trenches (e.g., Levitt & Sandwell, 1995; McAdoo & Martin, 1984), fracture zones in the Pacific ( $T_e$  described by the depth to the 600°–700°C isotherms; Wessel & Haxby, 1990) and the New England seamounts ( $T_e$  correlated with an isotherm in the range of 750–850°C; Zheng & Arkani-Hamed, 2002).

Watts et al. (2006) tentatively assigned a tectonic setting to each  $T_e$  estimate. They assigned seamounts in the range  $0 < T_e < 12$  km an “on-ridge” setting, seamounts with intermediate estimates  $12 < T_e < 20$  km a “flank ridge” setting, and seamounts with  $T_e > 20$  km an “off-ridge” setting. Considering the  $T_e$  estimates for the five hotspots analyzed in the present study, this classification is consistent with the “on ridge” setting of Azores, “flank ridge” of Great Meteor, and “off ridge setting” of Cape Verde, Madeira and Canary hotspots.

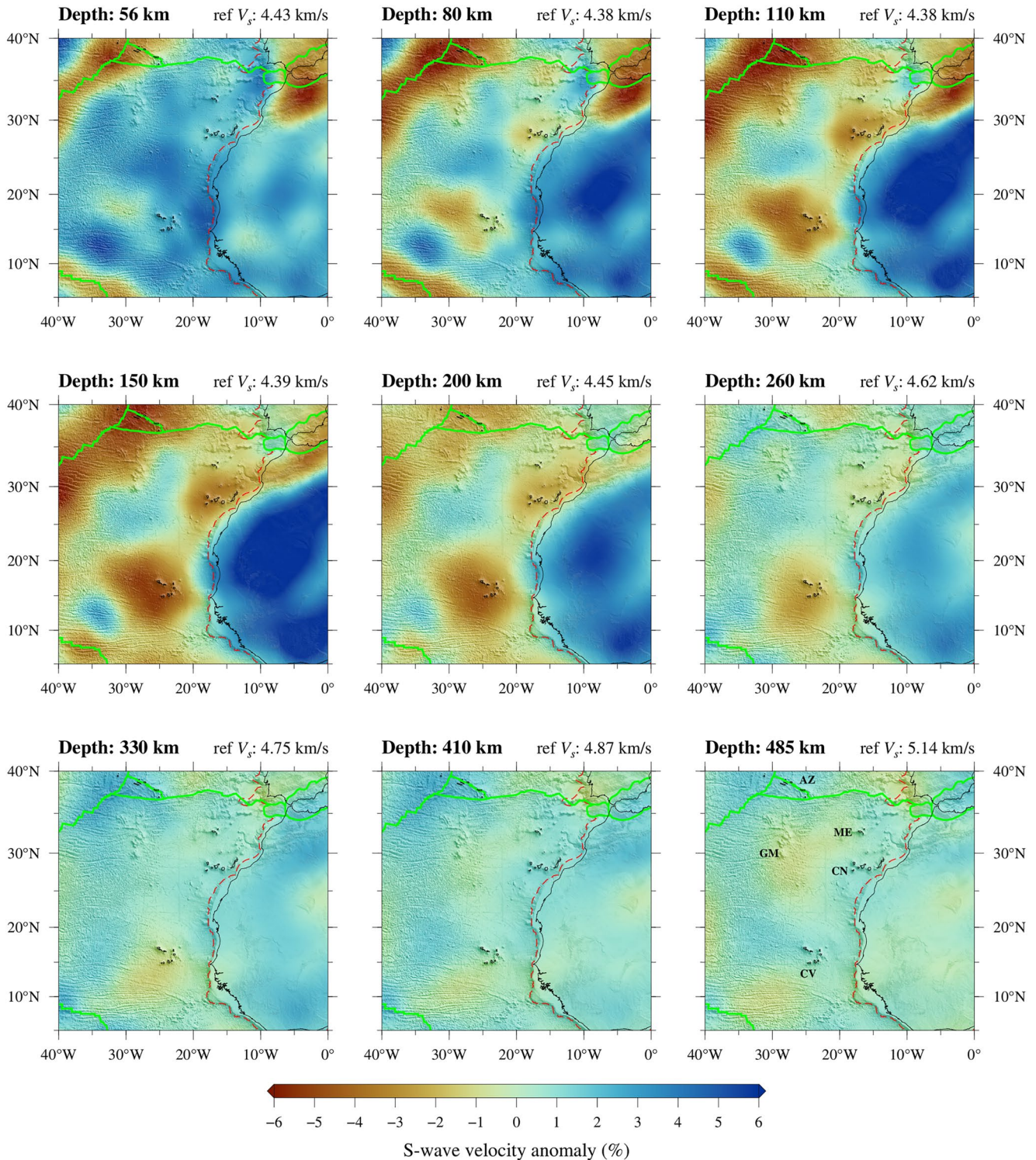
### 5.3. Comparison With Seismic Structure of the Upper Mantle

Compared to previous studies focused on specific areas, the regional approach adopted here enables a direct comparison between the lithospheric strength and the seismic structure of the lithosphere and upper mantle. Figure 9 shows the shear-wave velocity anomalies at different depths in the upper mantle from the surface-wave tomography model SA2019 (Celli, Lebedev, Schaeffer, Ravenna, & Gaina, 2020). At shallow depths (above 80 km), all the hotspots except Azores show high seismic velocities (bluish colors). At larger depths, all the hotspots show slow velocity values (reddish colors), with Azores and Cape Verde being slower than Canary, and Madeira being located at the periphery of the Canary anomaly. At 200 km depth, the negative velocity anomaly beneath Azores starts vanishing and Cape Verde exhibits the most negative anomaly. At 330 km, Azores presents a strong positive anomaly, Canary/Madeira neutral to slightly positive and the velocity model below Cape Verde remains negative. On the basis of the interpretation of these anomalies as temperature-derived, Azores is a “shallow” (source depth  $<200$  km) hotspot on top of a ridge, while Cape Verde is a “deep” hotspot characterized by a negative velocity anomaly in most of the upper mantle.

The comparison between Figures 5 and 9 reveals that, except for Azores, there is no correlation between the  $T_e$  estimated at the ECA hotspots and the amplitude of low shear-wave velocity anomalies in the underlying mantle. If the mechanical thickness of the hotspot were controlled by the presence of anomalously warm upper mantle material (on the basis of the interpretation of temperature-related seismic anomalies), one should expect similar  $T_e$  values for Cape Verde and Canary (possibly lower in Cape Verde) and a higher  $T_e$  in Madeira. Instead, our results show similar  $T_e$  values for Cape Verde and Madeira (Figures 5 and 8), and much higher  $T_e$  values beneath the Canary region. This lack of correlation between the elastic thickness and upper mantle temperature is in agreement with recent convection simulations of plume-lithosphere interaction at the Hawaiian Islands by Bellas et al. (2020), who showed that thermal erosion from a plume only perturbs the deepest part of the lithosphere, while the rest of the plate remains strong. The recent convection modeling of the interaction between an upper mantle plume and the lithosphere beneath the Canary Islands region by Negrodo et al. (2022) is also consistent with an almost negligible thermal erosion or viscosity reduction caused by the highly mobile mantle plume. The joint interpretation of the  $T_e$  and the seismic velocity structure performed in the present study permits extending this inference of the negligible effect of upper mantle temperature anomalies on the flexure to the ECA region.

### 5.4. On the Strength of the Lithosphere in the Canary Islands Region

Our results indicate the standard elastic thickness of the oceanic lithosphere beneath the ECA hotspots. However, the Canary hotspot is the notable exception, with a complex pattern of  $T_e$ . Our  $T_e$  results over the Canary Islands region are the highest values for intraplate volcanic islands worldwide (see e.g., Watts, 2001; Watts et al., 2006; Watts & Zhong, 2000); similar oceanic  $T_e > 50$  km values are found only along some trenches (e.g., Hunter & Watts, 2016; Levitt & Sandwell, 1995; McAdoo & Martin, 1984; Yang & Fu, 2018). The very high strength of the lithosphere at the Canary hotspot is apparently in contrast with the prominent negative shear-wave velocity



**Figure 9.** Seismic shear wave velocity anomaly at different depths in the lithosphere and upper mantle from the tomography model SA2019 (Celli, Lebedev, Schaeffer, Ravenna, & Gaina, 2020). Abbreviations are the same as in Figure 4.

anomaly down to 200 km (Figure 9). Fullea et al. (2015) suggested a sublithospheric temperature anomaly of 100 K on the basis of integrated geophysical and petrologic modeling. However, these authors did not find a significant thermal thinning or mechanical weakening of the lithosphere that could be related to the anomalously hot sublithospheric mantle. The convection modeling by Negredo et al. (2022) also showed negligible thermal

erosion of the lithosphere beneath the Canary region. This absence of weakening is in agreement with the high  $T_e$  estimates independently derived by the present work.

On the other hand, we note that the high  $T_e$  over the Canary Islands region is the most major discrepancy between this and previous studies, which yield  $T_e$  values ranging between 15 and 40 km (Canales & Dañoibeitia, 1998; Collier & Watts, 2001; Dañoibeitia et al., 1994; Kalnins, 2011; Llanes, 2006; Watts, 1994; Watts et al., 1997). These differences can be attributed to the use of different approaches and methods, different values of parameters, and/or different sizes of the study area. Focusing on the spectral estimates of  $T_e$ , it is important to note that the wavelet method provides a coherence estimate at every point on the data grid (see above). This allows us to map out the spatial variations in  $T_e$  over the Canary region instead of obtaining a single estimate of  $T_e$  within a finite-size window, representative for the entire analyzed area. To evaluate this effect, we calculated the spatial variations in  $T_e$  over the Canary region in the same way as in our main results as well as averaging over the entire area (see Figure S8 in Supporting Information S1). The coherence calculated averaging over the entire area yields a  $T_e$  value of  $26.7 \pm 2.8$  km for the Canary region, which is in good agreement with previous estimates.

A factor that can be related to the extremely high values of  $T_e$  in the Canary Islands region (Figure 5) is the presence of chemically depleted (therefore dry and more viscous) lithospheric mantle material (Neumann et al., 1995, 2002; Wulff-Pedersen et al., 1996). The lithospheric mantle xenoliths from the Canary Islands show ultra-refractory compositions that have been interpreted as residues after partial melting beyond the stability of primary clinopyroxene (Simon et al., 2008). Whether this hypothetical compositional strengthening can be applied to other hotspots or not needs further investigation. For example, Ito et al. (1999) obtained a significant viscosity increase due to water extraction from the mantle due to partial melting of the Iceland mantle plume.

The particular mechanical signature of the Canary Islands region could be affected by its close location to the edge of the West African Craton (Figure 9), and the possibility of the occurrence of edge-driven convection (EDC). This mechanism would generate small-scale convection caused by the contrast between thick cratonic lithosphere and normal or thin lithosphere and has been suggested by King and Ritsema (2000) as the origin of African plate hotspot volcanism. Numerical modeling of EDC in the Atlas Mountains by Kaislaniemi and van Hunen (2014) shows that the upwelling mantle in the convection cell makes decompression melting possible. After melt removal, a layer of dry and depleted material is left behind. This viscous (drier) and buoyant layer tends to stick to the base of the lithosphere, possibly increasing its strength. Recent numerical modeling efforts consistently find that the EDC mechanism is not consistent with the extent of a melting anomaly beneath the Canary Archipelago (Negredo et al., 2022) or even with the melting occurrence itself (Manjón-Córdoba Cabeza & Ballmer, 2021). Negredo et al. (2022) found that the interplay between an upper mantle thermal plume and the strongly heterogeneous lithosphere results in an asymmetric mantle flow and in lateral migration (away from the craton edge) of the plume head, in agreement with the overall westward decrease in the age of the islands. In turn, this laterally migrating pattern would prevent effective thermal weakening in the overlying lithosphere and  $T_e$  reduction.

As indicated by the complex SRC pattern with two rollovers (Figure 7), a lower  $T_e$  value would also be plausible. The complex  $T_e$  pattern can be understood in terms of the very wide range of volcanic ages in the Canary volcanic province. For example, the Bisabuelas (Grand-grandmothers, as termed by van den Bogaard, 2013) a small group of seamounts located 70 km southwest of El Hierro has ages up to 142 Ma, in contrast with the <1.1 Ma age of El Hierro island. Actually, van den Bogaard (2013) nominated the Canary Island Seamount Province as the most long-lived hotspot preserved on earth. We hypothesize that lower  $T_e$  values are likely reflecting load compensation of small seamounts located to the southwest of the archipelago, which show ages varying from 142 to 91 Ma (van den Bogaard, 2013) and therefore erupted in a much younger (<60 Ma) oceanic lithosphere than the Canary Islands. However, we cannot preclude that this complex SRC pattern is associated with the low shear-wave velocity anomalies in the underlying mantle.

We recognize that the complex  $T_e$  pattern in the Canary region is a complex issue that deserves further analysis. We hope to address this issue in the future by exploring the effects of multiple loads through synthetic modeling.

## 6. Conclusions

In this work, we investigate the mechanical structure of the lithosphere beneath the ECA hotspots by calculating their effective elastic thicknesses from an analysis of the Bouguer coherence calculated using wavelet transforms.

We find that the Azores hotspot is characterized by a  $T_e < 10$  km, whereas the Great Meteor, Cape Verde, and Madeira hotspots have intermediate  $T_e$  (15–30 km) values. In contrast, the Canary Island region is characterized by a much higher  $T_e$  (>50 km), forming the largest and most noticeable mechanical feature in the ECA.

We analyze the relation of the regional  $T_e$  distribution versus the age of the lithosphere and the temperature distribution on the basis of the plate cooling model. Averaged  $T_e$  values show the expected systematic increase with seafloor age and correlate with the depth of the 350–400°C isotherms, with an envelope of approximately 200–500°C. However, this trend is lost for oceanic lithosphere older than 130 Ma, which exhibits a  $T_e$  reduction.

When compared to the average values of the same seafloor age and with other regions in the world, we can conclude that Azores, Great Meteor, Cape Verde and Madeira have standard  $T_e$  values. In contrast, the very high  $T_e$  estimated for the Canary hotspot is found only in the oceanic lithosphere in some trenches. The comparison between the  $T_e$  distribution and the upper mantle seismic structure reveals the lack of correlation between the  $T_e$  estimated at the ECA hotspots (with the exception of Azores) and the amplitude of low shear-wave velocity anomalies in the underlying mantle. This comparison suggests a negligible effect of upper mantle temperature anomalies on the flexure of the ECA region.

The higher strength obtained for the Canary hotspot may be related to the highly depleted mantle composition in the area after melt extraction in a process that has not left a significant thermal imprint in the lithosphere. The complex  $T_e$  pattern in the Canary volcanic province may be the result of the very long, from the Late Jurassic to Recent, magmatic history of the region.

### Data Availability Statement

The GEBCO\_2022 Grid is available at [https://www.gebco.net/data\\_and\\_products/gridded\\_bathymetry\\_data/gebco\\_2022/](https://www.gebco.net/data_and_products/gridded_bathymetry_data/gebco_2022/). The EGM2008 gravity model is available at <https://bgi.obs-mip.fr/data-products/grids-and-models/egm2008-global-model/>. The global crustal model CRUST1.0 is available at <https://igppweb.ucsd.edu/~gabi/crust1.html>. The global model of oceanic crustal age is available at <https://www.earthbyte.org/category/resources/data-models/>. The global ocean sediment thickness model GlobSed is available at <https://ngdc.noaa.gov/mgg/sedthick/index.html>. The surface-wave tomography model SA2019 (Celli, Lebedev, Schaeffer, Ravenna, & Gaina, 2020) is available at <https://nlscelli.wixsite.com/nlseismology/sa2019>. Calculations were performed using the isowave software suite (Kirby and Swain). All figures were generated using the Generic Mapping Tools (GMT) package (Wessel et al., 2019). Several maps were represented using the Scientific Colour Maps, a suite of scientific, colour-vision deficiency friendly and perceptually uniform colour maps (Crameri, 2018; Crameri et al., 2020). The effective elastic thickness results are available at Jiménez-Díaz et al. (2023).

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