

2 **Recent Advances in Combined Positron Emission** 3 **Tomography and Magnetic Resonance Imaging**

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10 **ABSTRACT:** Hybrid imaging modalities combine two or more medical imaging techniques offering
11 exciting new possibilities to image the structure, function and biochemistry of the human body in
12 far greater detail than has previously been possible to improve patient diagnosis. In this context,
13 simultaneous Positron Emission Tomography and Magnetic Resonance (PET/MR) imaging offers
14 great complementary information, but it also poses challenges from the point of view of hardware
15 and software compatibility. The PET signal may interfere with the MR magnetic field and vice-
16 versa, posing several challenges and constrains in the PET instrumentation for PET/MR systems.
17 Additionally, anatomical maps are needed to properly apply attenuation and scatter corrections to the
18 resulting reconstructed PET images, as well motion estimates to minimize the effects of movement
19 throughout the acquisition. In this review, we summarize the instrumentation implemented in
20 modern PET scanners to overcome these limitations, describing the historical development of hybrid
21 PET/MR scanners. We pay special attention to the methods used in PET to achieve attenuation,
22 scatter and motion correction when it is combined with MR, and how both imaging modalities may
23 be combined in PET image reconstruction algorithms.

24 **KEYWORDS:** Hybrid Imaging, Medical Imaging, Multi-Modality Systems, PET, MR

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45 1 Introduction

46 Diagnostic and research medical imaging are intrinsically multimodal. Each modality offers specific
47 structural or functional information about the anatomical region being studied and/or treated, and
48 this information is associated with potential disease, injury, or the efficacy of a therapy. For example,
49 Computerized Tomography (CT) quantifies electron density, Positron Emission Tomography (PET)
50 determines concentration and distribution of a radiotracer in the patient, while Magnetic Resonance
51 (MR) imaging displays proton density and magnetic properties [1].

52 Imaging modalities have been evolving to constantly adapt to changing needs. New scanners
53 have benefited from technological advancements increasing their acquisition speed, spatial resolu-
54 tion, and image contrast. They now produce better images that allow physicians to better diagnose
55 diseases and carry out medical procedures with greater confidence. At the same time, medical
56 imaging researchers have worked on improving and combining the various equipment to obtain
57 complementary information or address the limitations of each of the imaging modalities. Obtaining

58 multimodal images is not straightforward. Moving the patient from one scanner to another in a
59 clinical setting is not easy and, depending on the position of the patient in the different scanners,
60 the acquired images could present a different spatial distortion. Thus, the need for precisely aligned
61 anatomical and functional imaging led to the development of the hybrid PET/CT scanner [2].
62 PET/CT scanners soon became widespread clinical equipment [3], while the development of com-
63 bined PET/MR scanners took longer because of various technical challenges of both technologies
64 [4, 5]. Although the combination of PET and MR was already mentioned in a PhD thesis in 1991
65 [6], and simultaneous PET/MR acquisition was demonstrated in 1997 [7], it took an additional 13
66 years for clinical systems to be commercially available [8].

67 Currently, several studies are demonstrating the benefits of hybrid PET/MR imaging compared
68 to PET/CT, including simultaneous acquisition to accurately quantify tissue and/or organ function-
69 ality from complementary modalities, reduced exposure to radiation and contrast-enhanced soft
70 tissue [9–18]. These benefits are the result of truly simultaneous data acquisition avoiding the need
71 for patient-specific CT images -and thus reducing the ionizing radiation dose- and the MR superior
72 soft tissue contrast allowing for better anatomical localization of activity in PET.

73 However, these new hybrid PET/MR scanners also come with new practical limitations for
74 clinical application compared to equivalent PET/CT systems, e.g., increased noise, different ar-
75 tifacts, attenuation correction (AC) issues, and new safety concerns. Moreover, implementing
76 both technologies in the same bore is challenging, as they caused interference with one another,
77 compromising the final image quality.

78 **2 PET instrumentation for PET/MR systems**

79 The typical detectors used in PET have been adopted from traditional nuclear physics to optimize
80 gamma photons detection with high sensitivity, fast response, high spatial resolution, and reasonable
81 energy resolution [19]. Good efficiency is needed to obtain an equivalent image signal-to-noise
82 ratio (SNR) with a reduced delivered dose to the patient. The total efficiency is a combination
83 of geometry (solid angle of the patient covered by the scanner) and detector intrinsic efficiency.
84 Nowadays, total-body scanners up to 1 or 2 m long have been introduced in order to increase scanner
85 efficiency [20]. High temporal resolution with a narrower coincidence time window aids to reduce
86 random coincidences (figure 1). Furthermore, below 1 ns, the coincidence time resolution (CRT)
87 enables the use of time-of-flight (TOF) information to reduce the uncertainty in the coincidence
88 positioning (figure 2), increasing the effective efficiency [21–23]. Equivalently, it is interesting to
89 precisely locate the interaction point of the photon within the detector to define narrower lines of
90 response (LORs) and reduce the depth of interaction uncertainty (DOI, figure 1), increasing image
91 resolution. Energy resolution allows distinguishing 511 keV photons from scattered photons (figure
92 1), which may achieve up to 50% of the coincidences in clinical scenarios [24, 25].

93 The strong magnetic fields used in MR, of the order of a few Tesla, may interfere with PET
94 detectors distorting the PET signal. In this section, we focus on the main components of the PET
95 detector and their compatibility with MR scanners:

- 96 • The scintillator crystal absorbs the gamma radiation and transforms it into a secondary light
97 signal.

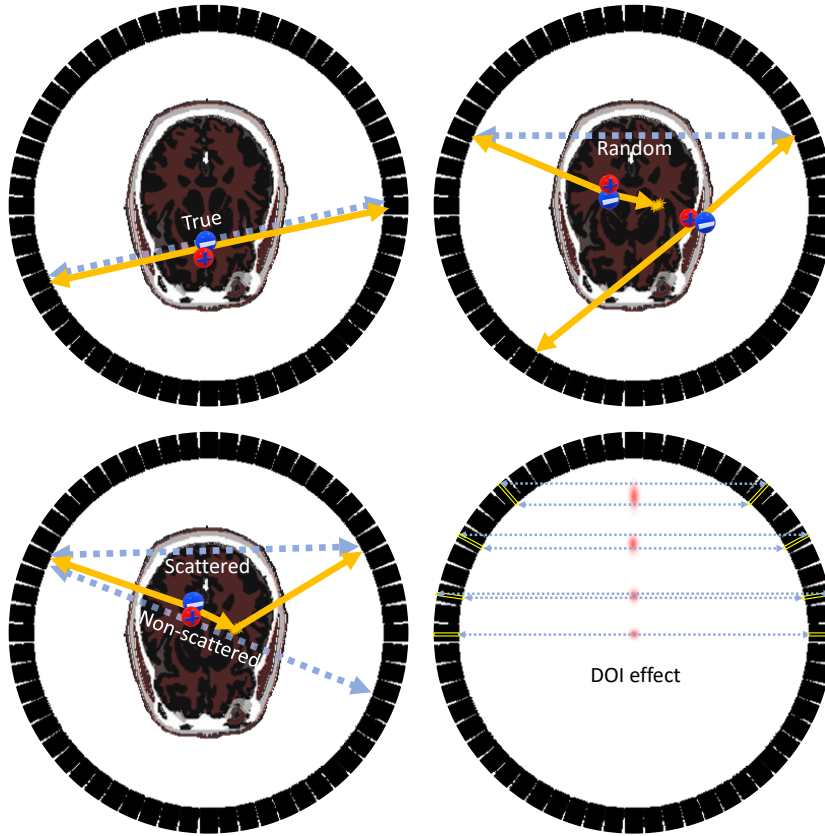


Figure 1. Graphical example of true (top-left), random (top-right) and scattered coincidences (bottom-left), and radial resolution loss of a point source caused by DOI effect (bottom-right).

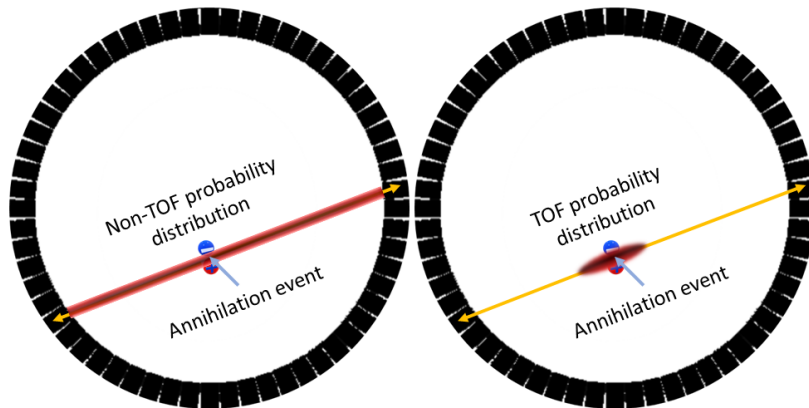


Figure 2. Scheme of the TOF information effect. Including TOF information (right) in the reconstruction process translates the uniform probability of an event occurring somewhere along the LOR to a Gaussian probability with full width at half maximum (FWHM) given by the CRT between both detectors.

- The photosensor, optically coupled to the scintillator crystal, generates an electronic pulse when a gamma photon is detected in the scintillator.
- The electronic circuit processes the electronic signal to obtain physical information about the detected gamma rays, and identify the coincidences needed for PET imaging.

2.1 Scintillation crystals

The scintillation detectors match PET needs with a reasonable trade-off between efficiency, energy resolution, and time resolution, hence they have been widely implemented. In a scintillation material, any gamma photon interaction results in an energy-proportional emission of photons in the visible energy range. Currently, there are different types of scintillation crystals (organic and inorganic in solid, liquid, and gaseous forms) [26]. In PET scanners, we commonly find solid inorganic crystals because of their high stopping power (defined as the energy lost per unit length) and high light response [27, 28]. In the context of PET/MR, scintillation materials do not present any contraindication.

The intrinsic efficiency of the scintillator is a function of its effective atomic number (averaged over its composition), its density, and the crystal thickness. The energy resolution is closely related to the photons emitted and their transport through the crystal. The number of photons emitted is quantitatively measured as the light yield, and the transport of photons is dominated by the refractive index of the material (the lowest, the better). The emission spectrum of the scintillator should match the photosensor sensitivity for optimal performance. It is also important to cover the crystal with reflectors to guarantee that a high fraction of the light can be collected. The time response in the scintillator has two main components: a fast-rising signal and a slow-falling signal (decay constant) that determines the time resolution [29]. In PET detectors, the scintillation crystals are usually pixellated to improve the spatial resolution of the photon interaction point. Some manufacturers have developed monolithic scintillators too [30–32].

Lutetium-based scintillators LSO (Lu_2SiO_5) [33] and LYSO ($\text{Lu}_{2(1-x)}\text{Y}_{2x}\text{SiO}_5$) have become the trend in modern PET scanners because of their high stopping power, high light yield, and fast time response [28]. Other proposed materials are GSO (Gd_2SiO_5) [34, 35], since it has a low price and acceptable performance, making it suitable for long axial field-of-view (FOV) scanners, and GAGG ($\text{Gd}_3(\text{Ga}, \text{Al})_5\text{O}_{12}$) or LaBr_3 [28], because of their superior performance. $\text{NaI}(\text{Tl})$ was the first scintillation crystal used in PET [36], but it was soon substituted by other materials with higher stopping power, such as BGO ($\text{Bi}_4\text{Ge}_3\text{O}_{12}$). In contrast, BGO has a low light output yield compared with $\text{NaI}(\text{Tl})$ and other scintillators [19]. In recent times, BGO has drawn the attention of researchers since it offers the possibility of using Cherenkov radiation to improve the coincidence time resolution in TOF-PET scanners [37, 38]. In table 1 we summarize the main properties of these materials. Other authors also proposed cheaper detectors with lower sensitivity and high coincidence time resolution based on Cherenkov radiation, such as plastic scintillators [39] or PbF_2 [40].

2.2 Photosensors

The photosensor converts the optical response of the scintillator to an electrical signal that is later processed. To perform energy spectroscopy of the incident gamma rays, this device must

Table 1. Main properties of common PET scintillators. Adapted from Yu et al. [28] and Lewellen et al. [19]. The refractive index for GAGG was taken from [41].

	NaI (TI)	BGO	GSO	LSO	LYSO	LaBr₃	GAGG
Chemical Formula	NaI	Bi ₄ Ge ₃ O ₁₂	Gd ₂ SiO ₅	Lu ₂ SiO ₅	Lu _{2(1-x)} Y _{2x} SiO ₅	LaBr ₃	Gd ₃ (Ga, Al) ₅ O ₁₂
Z_{eff}	51	74	59	66	60	47	48
Density (g/cm³)	3.67	7.13	4.89	7.4	7.2	5.3	6.63
Light yield (ph/keV)	41	9	10	31	30	67	54
Wavelength (nm)	410	480	440	420	420	370	540
Decay constant (ns)	230	300	60	40	41	25	94
Refractive index	1.85	2.15	1.85	1.82	1.81	1.95	1.90

138 maintain the proportionality between the light generated and the electrical pulse. Spatially sensitive
 139 photosensors are also required to provide the gamma interaction point in the detector [19]. Another
 140 quality required for the photosensors in recent multimodal PET/MR scanners is being able to operate
 141 under magnetic fields. The two most common photosensors employed in PET are photomultiplier
 142 tubes (PMTs) and silicon photomultipliers (SiPMs).

143 PMTs are made of a vacuum tube that encloses three main parts: the photocathode, a
 144 sequence of dynodes, and the anode [26]. When the photons hit the photocathode, the electron
 145 response generated is amplified in the dynodes sequence, obtaining a signal in the anode proportional
 146 to the energy deposited in the scintillator. The PMT has been the reference photosensor for PET
 147 scanners since the beginning of the technique, but its disadvantages are encouraging the use of other
 148 photosensors in the field. Some of them are its sensitivity to magnetic fields [42], requiring special
 149 shieldings to make them compatible with MR [43], its bulky size, and its expensive cost because of
 150 their complex manufacturing process.

151 SiPMs have overcome most of the disadvantages of PMTs: they have similar gain, fast response,
 152 low manufacture expenses, possible one-to-one coupling with the scintillation pixels, and especially,
 153 compatibility with MR scanners [44]. A few drawbacks of SiPMs are non-linear response or lower
 154 photon detection efficiency. Although SiPMs add noise due to thermal factors, which should be
 155 considered when working together with MR, with correct temperature controls, SiPMs are the best
 156 option to replace PMTs in combined PET/MR scanners. SiPMs are conformed of a matrix of
 157 avalanche photodiodes (APDs). The APDs are semiconductor detectors that create an avalanche of
 158 hole-electron pairs when a visible or ultraviolet photon is absorbed in a depleted region. The APD
 159 itself can be used as a photosensor, but they suffer from high noise in the amplification process,
 160 compromising the temporal resolution. For this reason, in SiPMs, the matrix of APDs works on
 161 Geiger mode, in which they detect single photons when a signal is generated over a threshold [45].
 162 The multiple cells can provide a proportional signal to the input light by summing all the activated
 163 cells. Currently, SiPMs have displaced PMTs in most of the modern TOF-PET systems [46–51]
 164 and combined PET/MR scanners [52–56]. However, there is still space for other MR-compatible
 165 photosensors, such as APDs [57, 58], being an active research field.

166 **2.3 Electronics**

167 When a photon interacts with the detector and deposits an amount of energy above a given threshold,
168 the electronics will detect the event (single event) and from the electronic pulse generated in the
169 photodetector, the electronic circuit behind will obtain information about the time stamp, deposited
170 energy, and location of the gamma photon interaction [26]. To avoid time-energy walk effects,
171 the time stamp can be obtained analogically with the help of a constant fraction discriminator
172 (CFD) that estimates the moment in which the pulse exceeds a constant fraction of its maximum
173 amplitude. The total charge accumulated in the photosensor is used to estimate the energy of
174 the photon. Usually, there are low- and high-energy discriminators to discard events far from the
175 511 keV peak. The gamma interaction point is estimated by weighting the multiple outputs of a
176 position-sensitive photosensor.

177 First PET detectors used a quadrant-sharing scheme, where up to four PMTs would read the
178 light produced in an array of scintillators, and the measurement of the light shared by each PMT
179 allows to place the gamma interaction in the XY plane of the detector [59]. In order to reduce the
180 number of electronic signals to be read, these were multiplexed by means of resistor networks. Most
181 modern PET scanners today, however, used a separate SiPM photosensor for each small scintillator
182 element, as it is the case in the Siemens Biograph Vision PET/CT scanner (Siemens Healthineers,
183 Erlangen, Germany) [60], and read every signal from each separated photosensor. This increases
184 the light received by each individual photosensor and allows for instance for improvements of the
185 timing information and energy resolution of the newest scanners. This comes at the price of a
186 many-fold increase in the number of signals that need to be read, decoded, and digitized. Dedicated
187 ASIC electronics are at the heart of most modern PET scanners, in order to keep the cost of the
188 electronics under control.

189 Recalling the interest of detecting pairs of annihilation photons, single events are processed,
190 either by a hardware coincidence unit, or by a software coincidence algorithm, to identify pairs of
191 events within a time window of the order of a few nanoseconds. This is what we call electronic
192 collimation, and we call a coincidence event to every pair of single events within the time window.
193 The TOF is very useful to improve the decay location within the LOR, and the CRT has been
194 continuously improving over the last years, being now close to 200 ps in commercial clinical
195 scanners [61].

196 **3 PET/MR configurations**

197 Since the first PET/MR prototype was developed in 1997, several approaches have been considered
198 to acquire PET and MR images in a single scanner [62, 63], although their overall designs are
199 substantially different. They are normally categorized into sequential or simultaneous systems, the
200 last category including insert-based and integrated PET/MR systems.

201 **3.1 Sequential PET/MR**

202 The first MR-compatible PET systems used PMTs coupled to scintillating crystals to detect gamma
203 rays and amplify the signal. However, PMTs are overly sensitive to the static fields used in MR,
204 making the data compress on one side of the detector [64]. The solution was to locate the PMTs

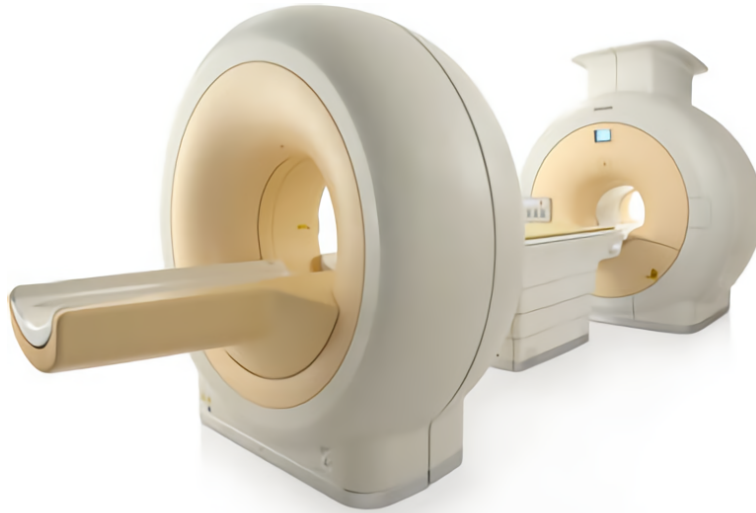


Figure 3. Commercially available Philips TF Ingenuity sequential PET/MR scanner: PET and MR scanners placed in the same room but maintaining enough distance to avoid electromagnetic interference. Image courtesy of Philips Healthcare.

205 outside the magnet and couple them to the scintillating crystals placed inside the scanner by using
206 long optical fibers. This structure was used to simultaneously acquire for the first time PET
207 images and ^{31}P NMR spectra from isolated, perfused hearts and performing *in vivo* rat brain
208 studies [65]. Multiple improvements were proposed afterward, but these evolutions demonstrated
209 limited performance and were used only for pre-clinical imaging studies in small animals [66].
210 Actually, the first PMT-based human PET/MR scanner -the Philips TF Ingenuity (Philips Healthcare,
211 Amsterdam, Netherlands)- allowed only sequential imaging. This scanner was composed of PET and
212 MR scanners placed in the same room but maintaining enough distance to avoid electromagnetic
213 interference (Figure 3). Data were acquired from each modality independently and posteriorly
214 fused using software in a post-processing step. Nevertheless, this approach was not able to solve
215 misalignment issues with moving organs. The first two clinical whole-body PET/MR systems were
216 installed in 2010 [8].

217 3.2 PET/MR inserts

218 Early approaches for the advance of combined PET/MR scanners started by developing PET inserts
219 that were compatible with already existing high-field animal MR systems [67–70]. PET inserts are
220 basically a supplement to the MR scanners consisting of radially arranged LSO detector modules
221 that could be coupled inside the bore of the MR scanner. The development of APDs allowed these
222 PET detectors to be positioned in the bore of the MR scanner, and thus enabled simultaneous data
223 acquisition.

224 At the same time these advancements in the preclinical PET/MR scanners were taking place,
225 and with a focus on developing a completely integrated PET/MR scanner in the future, Siemens
226 first designed several prototypes of the so called BrainPET/MR scanner [71, 72]. The BrainPET
227 consisted of a PET insert which was placed into a standard, slightly modified 3T MAGNETOM

228 Trio MR scanner, and demonstrated its ability to simultaneously image a human brain in a one-bed
229 position with both PET and MR. The PET ring insert had an internal diameter of 36 cm in which
230 two dedicated MR head coils optimized for minimal PET attenuation were positioned. As these
231 scanners were never commercially available, only four BrainPET prototypes were installed around
232 the world between 2006 and 2010.

233 **3.3 Integrated PET/MR systems**

234 Currently Siemens, General Electric (GE Healthcare, Waukesha, Wisconsin, USA), and United
235 Imaging (United Imaging Healthcare, Shanghai, China) are the only vendors to offer a fully inte-
236 grated whole body and simultaneous acquisition PET/MR system. As the PET detector is integrated
237 within the MR gantry –between the gradient and the radiofrequency coils–, these hybrid scanners
238 allow for completely simultaneous PET and MR acquisitions. Contrary to what happens in simul-
239 taneous systems based on PET inserts, the PET detector is completely integrated and not noticeable
240 from outside the gantry, and does not use any space in the bore, maintaining a 60-cm wide diameter
241 bore (Figure 4). This full PET and MR integration allowed for the whole-body application of
242 PET/MR systems.

243 In this sense, Siemens designed the first PET APD detectors that permitted the full integration
244 of PET and MR while maintaining the performance of each modality, and introduced the first
245 fully integrated whole-body MR-PET scanner, called Biograph mMR [57]. The Siemens Biograph
246 mMR system received the CE mark and FDA approval in 2011 and was the first company to offer
247 simultaneous PET/MR acquisitions, with the first systems based on APDs installed in 2010 [73–
248 75]. The actual architecture design also included integrated cooling features to guarantee optimal
249 PET performance, as well as dedicated shielding to remove magnetic field interfering PET data.
250 Nonetheless, APDs temporal performance is slower than that of PMTs, therefore, limiting the use
251 of these detectors in TOF applications.

252 General Electric decided to focus its PET/MR developments on SiPM, which proved to be
253 successful in combined PET/MR applications [76]. Eventually, General Electric introduced the GE
254 SIGNA TOF PET/MR, which received its 510K clearance and CE mark in 2014, becoming the first
255 whole-body integrated PET/MR scanner with TOF capabilities in the market [77].

256 Later on, United Imaging joined the effort in developing combined PET/MR systems, integrat-
257 ing the strengths of next-generation SiPM-based detectors with TOF capabilities (480 ps timing
258 resolution) and a 3T MR in the uPMR 790 HD TOF PET/MR system. This scanner received FDA
259 clearance in 2019.

260 **4 PET image correction and reconstruction in PET/MR systems**

261 Annihilation photons in PET are subject to both attenuation and scatter as they travel through matter
262 [78–81]. Consequently, the amount of detected photons in each LOR is reduced. Additionally,
263 the photons can be scattered due to the Compton effect but still reach the PET detectors. Precise
264 corrections for both physical phenomena are required to produce quantitative images reflecting the
265 true spatial distribution of the radiotracer.



Figure 4. Integrated PET/MR Scanners. Commercially available integrated PET/MR scanners: Siemens Biograph mMR (top), GE SIGNA PET/MR (middle) and UI uPMR790 PET/MR (bottom). Images courtesy of Siemens Healthineers, General Electric Healthcare, and United Imaging Healthcare.

266 **4.1 Attenuation correction**

267 Loss of counts due to attenuation and scatter increases image noise, artifacts, and distortion. Figure
268 5 shows how, without AC, significant artifacts may occur on PET scans including:

- 269 • prominent activity at body surface edges due to relative lack of attenuation at the surfaces
270 compared to deeper structures.
- 271 • distorted appearance of areas of intense activity (e.g., urinary bladder) due to variable degrees
272 of attenuation in different directions of activity originating from these areas.
- 273 • diffuse and relatively increased activity in tissues of relatively low attenuation (e.g., lungs).

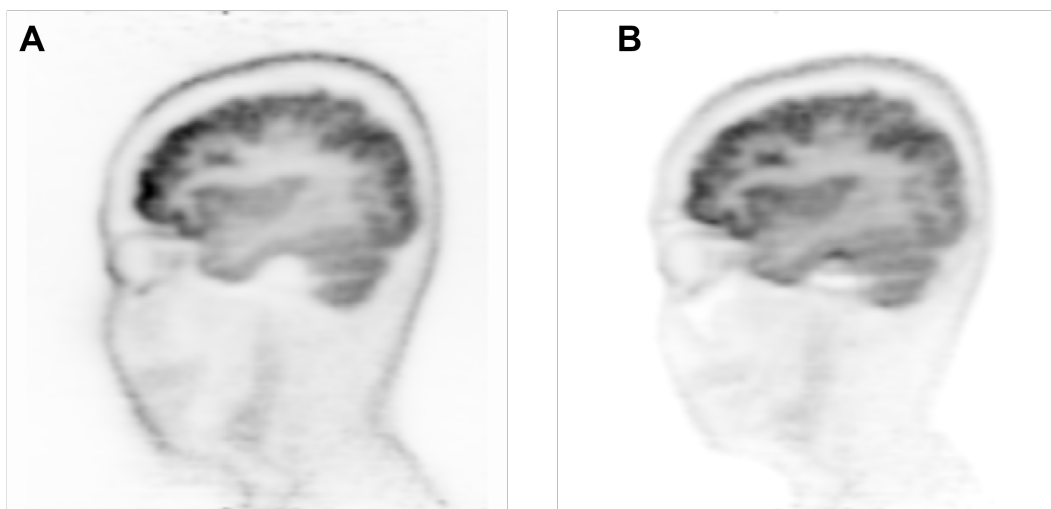


Figure 5. Sagittal slice of a ^{18}F -FDG PET study showing the reconstructed image without attenuation correction (A) and with attenuation correction (B). Note the background noise, the prominent uptake at the skin and the levels in the frontal lobe.

274 The measurement along each LOR can be corrected if the properties of the object are known
275 [82]. Nevertheless, constructing PET AC maps for use in simultaneous PET/MR systems is
276 challenging because no direct relation exists between linear attenuation coefficients (μ) and MR
277 signal intensity [83], contrary to what happens with AC maps and the intensity of CT images
278 [84]. Additionally, treating bone as soft tissue in MR-derived AC maps for PET/MR AC leads to a
279 substantial underestimation in the analysis of PET tracer distribution [85].

280 In the case of the fully integrated PET/MR scanners accessible in the market, AC maps are
281 usually generated from the MR data acquired using a “Dixon Volumetric Interpolated Breath-hold
282 Examination” (Dixon-VIBE) sequence or a “Liver Acquisition with Volume Acceleration-Flexible”
283 (LAVA-Flex) sequence. Typically, four classes of tissue are segmented (background, fat, lung and
284 soft tissue) based on the Dixon-VIBE or LAVA-Flex data, and discrete linear attenuation coefficients
285 are allocated to each of these classes. Recent developments added bone information over these maps
286 based on atlas. The main drawback of these methods rests in the inaccurate assignment of bone
287 as soft tissue in the four-class segmentation approach, leading to a significant bias in the PET

288 measurements. Similarly, assigning discrete and unvarying linear attenuation coefficients does not
289 describe the entire diversity of attenuations in actual tissues.

290 Hence, most of the early developments focused on trying to solve this issue in the head and
291 neck [86, 87], principally differing in the type of semantic representation –segmentation [88, 89]
292 or atlas-based [90, 91] methods– used to explain the resulting AC maps. In this setting, MR-based
293 pseudo-CT synthesis (estimation of the patient-specific CT from the acquired MR data) can assist
294 in the generation of continuous, more detailed and accurate MR-based AC maps, with a comparable
295 quality to those AC maps obtained from the patient-specific CT, and several groups worked towards
296 this direction.

297 Several different early developments started trying generating a head & neck pseudo-CT and
298 following diverse approaches such as image registration and pattern recognition techniques to match
299 MR and CT data [92], atlas-based methods grounded on non-rigid registration of a parametric atlas
300 [90], enhanced approaches based on non-rigid registration of a non-parametric multi-atlas followed
301 by a label fusion step depending on patch similarity measures [93], or the use of Convolutional
302 Neural Networks (CNNs) to synthesize pseudo-CT images [94], even using incomplete (head &
303 neck MR images and restricted local portions of CT) databases [95]. An extensive assessment of
304 several of these methods demonstrated that many are accurate enough for clinical practice [96].

305 Figure 6 illustrates a comparison among AC maps of the head and neck obtained from a real CT
306 (Figure 6A), by using a segmentation by classes approach (Figure 6B) and a pseudo-CT synthesized
307 employing the method described by Torrado-Carvajal et al. [91] (Figure 6C), in addition to their
308 matching attenuation corrected PET images for each of these approaches (Figures 6D-F) and fused
309 with an anatomical T1w MR image (Figure 6G-I). This figure shows how the PET image corrected
310 using the AC map derived from the segmented image demonstrates a decrease in signal in the most
311 internal side of the skull, along with frontal and occipital areas, because of the inaccurate estimation
312 of the skull, whereas the PET corrected using the AC map derived from the pseudo-CT is more
313 similar to that obtained with the AC map derived from the actual CT of the patient.

314 In view of these results, it seemed the MR-based AC challenge was solved in combined
315 PET/MR scanners. Nevertheless, translating these same approaches to whole-body applications
316 was not straightforward. The greater variability amongst subjects (i.e., morphology, body mass
317 index, pathology, sex) makes it problematic to create and establish an atlas that satisfactorily
318 describes the entire population. Additionally, the usage of the aforementioned approaches in other
319 regions different than the head and neck entails an amalgamation of rigid and non-rigid registration
320 (atlas deformation). Subsequently, most of these approaches present inadequate performance when
321 applied to whole body AC [81, 97, 98]. Thus, several research efforts have been made in recent
322 years to tackle this challenge.

323 For instance, preliminary efforts tackling with the pelvic region reported semiautomatic ap-
324 proaches relying on multiple MR sequences (Dixon and zero-echo-time (ZTE) images) to estimate
325 a pseudo-CT from the combination of the MR sequences [99]. Still, these approaches required
326 a manual correction stage throughout the bone segmentation, consequently restricting its clinical
327 applicability in real scenarios. As an alternative to classical atlas-based methods and the initial
328 semiautomatic solutions, and considering the growing awareness in deep learning (DL) approaches,
329 the use of CNNs for MR-based AC maps synthesis started to be studied for pelvic studies. Several
330 proposals focused on estimating multiparametric MR models to synthesize pseudo-CT images of

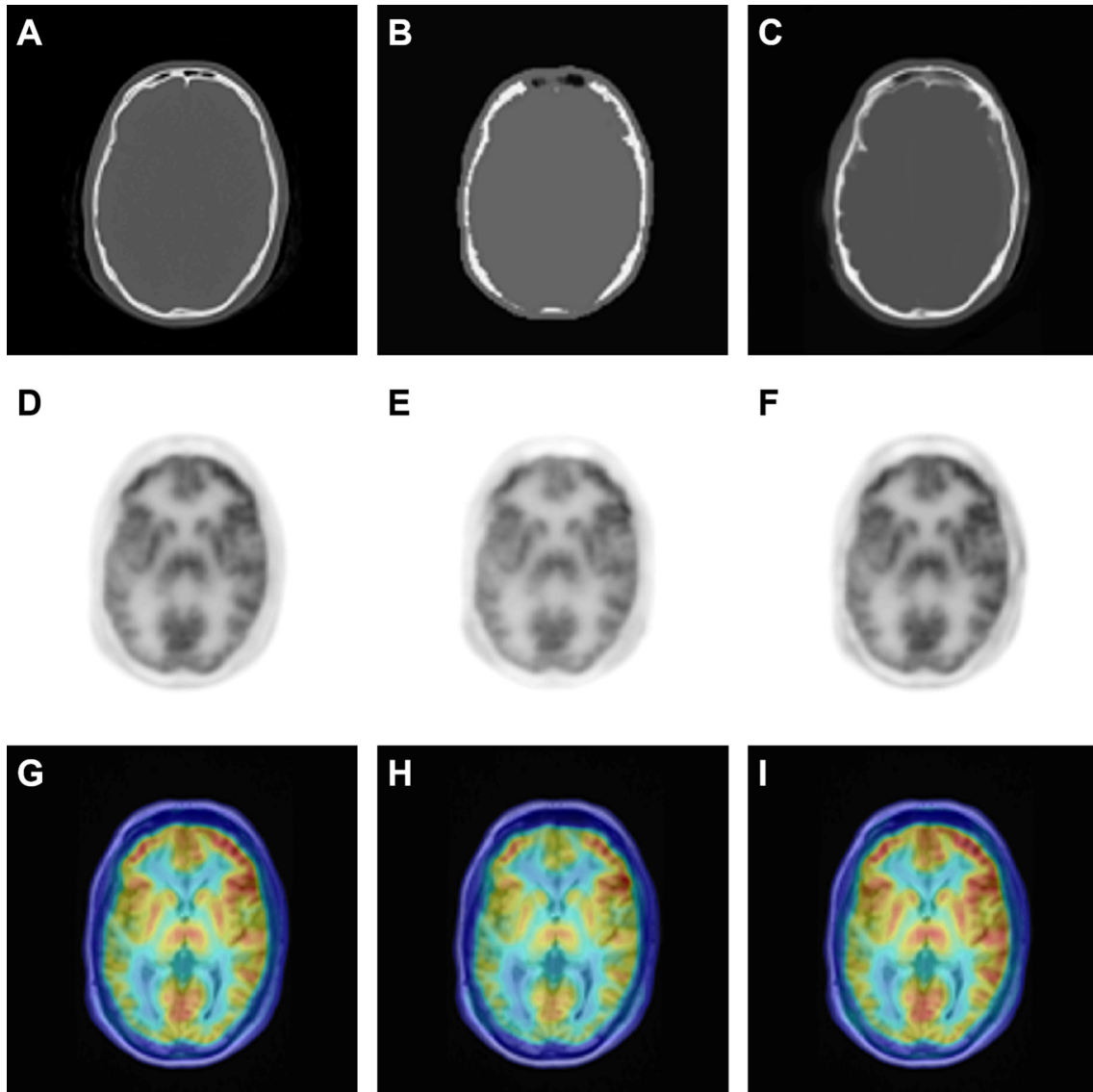


Figure 6. Axial slices of AC maps of the head and neck obtained from a real CT (A), and from MR images with a method based on segmentation by classes (B) and with a pseudo-CT synthesis method (C), as well as the 18F-FDG PET study images corrected with each of these methods (D-F) and the representation of each of these PET images fused with an anatomical MR T1 image of the patient (G-I). From Torrado-Carvajal et al. [81].

331 the pelvis based on Dixon images and additional ultra-short-echo-time or ZTE images [100]. While
 332 their preliminary results were promising, these approaches present the drawback of requiring dedi-
 333 cated MR sequences that involve added MR acquisition times. Nevertheless, the rapid advancement
 334 in CNN architectures and flexibility derived in novel developments that could map between the four
 335 MR Dixon-VIBE images (water, fat, phase and counterphase) and CT images, generating precise
 336 pseudo-CT images of the pelvis without the need for acquiring extra sequences [101, 102]. The
 337 main challenge of these approaches was that neglecting air pockets in the pelvis AC maps could

338 lead to inaccuracies in reconstructed PET images. However, these innovative methods based on
339 CNNs even showed the capability of predicting air pockets and include them in the AC maps [103].

340 Even the described AC methods presented some challenges when applied to chest imaging,
341 due to its high anatomical complexity and inherent cardiac and respiratory motion [104]. On the
342 one hand, most of the approaches neglected mapping the continuous distribution of attenuation
343 coefficients inside the lungs, which fluctuates both within a given subject and among subjects,
344 and depend on the breathing cycle, the lung inflation, gravitational dependence, and pathology.
345 Although MR was shown to be of help when estimating continuous lung coefficients in MR-based
346 AC [105]. On the other hand, respiratory and cardiac motion can lead to misalignments between the
347 AC map and PET, resulting in improper PET data quantification, for instance, in the quantification
348 of cardiac perfusion and metabolism [106, 107]. Deeper insight on these challenges will be covered
349 in the motion correction section.

350 There are other factors that should be mentioned in the context of MR-based AC. In integrated
351 systems, radiofrequency MR coils can also attenuate gamma ray photons when placed between the
352 subject and PET detectors. Equipment manufacturers provide attenuation maps for most coils so
353 that they can be automatically combined with the human maps in the AC procedure. Moreover,
354 while both arms should be considered in the AC map, they are truncated in the FOV in the MR
355 images, causing an underestimation of the tumour activity by 10 to 25% [108]. Nevertheless, the
356 truncated arms can be inferred using the maximum-likelihood reconstruction of attenuation and
357 activity procedure [109]. It should be also noted that all these methods and anatomical areas are
358 subject to distortions in the MR-based AC maps due to the presence of metallic implants. To answer
359 the issues derived from imaging patients with implants, MR images and AC maps can be improved
360 by using specialized sequences that allow imaging near metallic implants (MAVRIC or SEMAC
361 sequences), recovering the lost signal by means of image post-processing methods, or completing
362 the AC map based on joint estimation of emission and attenuation [110, 111].

363 4.2 Scatter correction

364 In PET, scatter typically denotes Compton scattering. In this process, one or both annihilation
365 photons interact with an electron, losing some energy and changing their direction [26]. If the
366 scattered photon is eventually detected, the resulting LOR will not correspond with the original
367 one (see figure 1), resulting in a smooth background and artifacts that reduces the contrast of the
368 images and may affect especially when many of these events are accumulated [112–114]. In clinical
369 acquisitions, the proportion of scattered events may be up to the 50% [24, 25, 115], making it a very
370 significant correction. Many modern scanners work with a narrow energy window around 511 keV
371 (usually from 400 to 650 keV [34, 57, 70, 72]) to discard most of the scatter events. Nevertheless,
372 scatter corrections are still required even with a narrow energy window.

373 Multiple approaches have been proposed in the literature to perform scatter corrections in
374 PET. Their main differences are the way in which the scatter is estimated. Once calculated, the
375 scatter contribution may be subtracted from the data or added to the forward projection model.
376 A few authors implemented experimental approaches based on multiple energy windows (spectral
377 approaches) [116–118]. The basic hypothesis is that the whole energy spectrum of the scattered
378 counts or at least the integral of that spectrum –from the lower energy discriminator to the upper
379 one of the photopeak window– could be estimated given a big enough number of events below

380 and above the photopeak. These approaches typically include supplementary information to obtain
381 potentially further accurate outcomes, aiming at withdrawing the scattered counts, voxel by voxel.

382 In the past, some authors proposed analytical methods based on fitting a Gaussian function to
383 the scatter tails at the edge of each projection (assuming this contribution was mainly scatter) [116].
384 This approach obtains good results in brain imaging, as the activity and the scattering medium is
385 homogeneously distributed and focused on the center of the FOV, resulting in a simple gradually
386 varying scatter distribution. However, this approach fails in whole body applications, as the scatter
387 tails obtainable for fitting are way shorter (since the body occupies a bigger part of the FOV) and
388 the scatter distribution comprises further structures.

389 The most commonly used method has been the single scatter simulation (SSS) [119–121]. SSS
390 includes a realistic model of scattering based on the activity emission map and the attenuation map
391 (either obtained from CT or MR) to integrate the single scatter contribution of all the voxels in
392 every LOR. Since scatter contribution is smooth, the uncertainties of MR-based attenuation maps
393 are way less relevant than for the attenuation correction. The SSS needs a global scale factor that
394 is usually obtained with tail fitting strategies based on the fact that no true activity is expected in
395 the edges of the FOV and only scatter counts may be present [122, 123]. However, tail fitting may
396 result in halo artifacts around high organ-to-background regions, such as the bladder or kidneys, or
397 create biased estimations if there are arms or other sources in the border of the FOV [113, 114].
398 Other intrinsic drawbacks are the missing detector modeling [124], or multi-scatter events that may
399 account from 15 to 40% of total scattered events [24, 25], resulting in wrong contrast values [125].
400 Some of these drawbacks were solved by Watson et al. in the most recent version of SSS [126],
401 adding double scatter simulation (DSS) and considering all the physics in the scanner to estimate
402 the scatter global scale.

403 Pure Monte Carlo (MC) simulations are considered the most accurate method for scatter
404 estimation [116, 124, 127, 128], and might become a substitute to overcome the drawbacks of SSS
405 thank to the higher computational power of current graphics processing units (GPUs) and clusters
406 of central processing units (CPUs) [124, 129]. To improve SSS robustness, some authors combined
407 MC global scaling factor from a quick simulation with the SSS scatter distribution [130, 131].
408 MC methods were combined with variance reduction techniques to increase the computational
409 performance [116, 128], but it was not until GPU parallelization that they were considered for
410 practical use [124, 127].

411 DL methods may help to speed up and enhance the scatter correction performance of the
412 aforementioned methods. A few authors relied on DL methods to achieve quick MC-equivalent
413 sinogram scatter estimations using the prompts and attenuation sinograms or SSS estimation as
414 input [132–135]. Long MC simulations are treated as the gold-standard output targeted by the
415 network in order to speed up the prediction, and the network estimation is incorporated in standard
416 reconstruction algorithms. Prats et al. proposed a different approach based on a gradient boosting
417 decision tree to tag true and scatter events [136], despite they did not achieve MC scatter fraction
418 values. Image post-processing approaches have been proposed as well, in which a non-corrected
419 PET image is transformed into a corrected PET image, usually joining attenuation and scatter
420 correction in the process [137–140]. These promising results provide a practical solution for
421 standalone PET scanners in the clinic. Nonetheless, DL methods are susceptible to outliers,
422 resulting in artifacts, quantification errors, missing lesions or false-positive results, so special

423 attention must be taken in these approaches [140].

424 4.3 Motion correction

425 Motion is difficult to avoid in long lasting PET studies, degrading the overall image quality and
426 offsetting the advantage of employing a high-resolution scanner. Typical motion sources come from
427 the inherent difficulty of the patient to stay still for long periods of time (voluntary or involuntary
428 movement), but more specifically from physiological processes such as respiratory and cardiac
429 motion. As the PET data take many minutes to acquire, the motion effects of the entire respiratory
430 cycles or cardiac cycles are averaged together in the final reconstructed image. Figure 7 shows how,
431 without motion correction, significant artifacts which may occur on whole-body PET. These effects
432 result in a:

- 433 • blurring of the final PET images and severe artifacts when motion has large amplitude,
434 resulting in spatial resolutions close to 1 cm.
- 435 • mismatch between the PET and MR data, which can even produce improper AC during PET
436 reconstruction, especially at soft tissue, bone, and lung interfaces.

437 As a potential solution, in combined PET/MR scanners, the simultaneously acquired MR data
438 can be employed to derive high-temporal-resolution motion estimates, removing the need for other
439 external tracking systems.

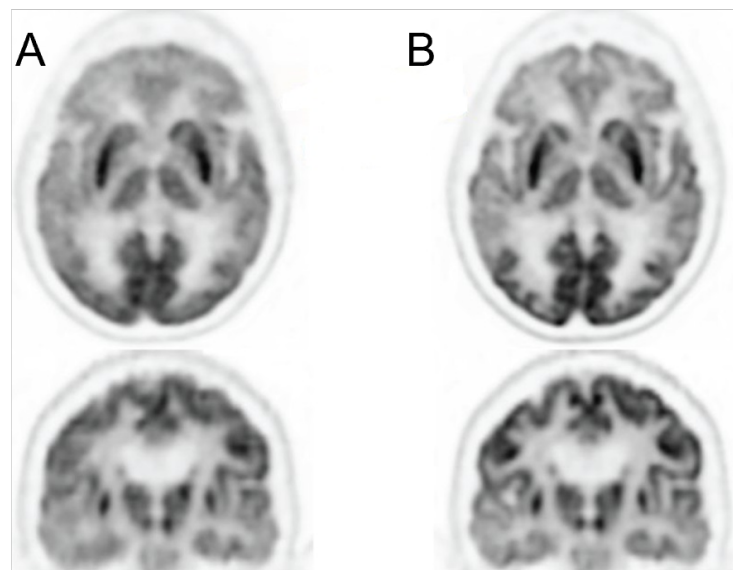


Figure 7. Axial (top) and coronal (bottom) slices of a ^{18}F -FDG PET study showing the reconstructed image without motion correction (A) and with rigid motion correction (B). Adapted from Spangler-Bickell et al. [141].

440 Numerous methods have been developed in the case of neurological PET studies [142, 143],
441 starting with simple methods to realign individual frames (shorter time reconstructions) to a fixed
442 reference position and summing all of them to generate the final single volume. However, these
443 methods neglect the availability of simultaneously acquired MR data that could be used to track
444 motion. In this sense, several approaches have been assessed in brain imaging, such as using the

445 information extracted from echo planar imaging (EPI) series or embedded cloverleaf navigators
446 (CLNs) [144], or the use of MR navigators sampled between other protocolled MR sequences
447 during simultaneous PET/MR brain scanning [145].

448 Recent studies have demonstrated how a failure to perform motion correction results in large
449 artifacts among high-motion subjects and how correcting for motion alleviates most of the bias, and
450 specifically correcting for inter-frame motion delivers most of the benefit [146].

451 Whole-body PET imaging, and in particular cardiac imaging, is restricted by the degradation
452 due to motion. Due to respiration, the diaphragm and the liver move in the range of 7–28 mm and
453 5–17 mm, respectively, and, during free breathing, the heart moves in the range of 4.9–9 mm due
454 to its own heartbeat and to the movement of the chest due to respiration. These motions can be
455 managed by using gating, which can remove most of the artifacts but at the cost of increased noise in
456 the final reconstructed images and, sometimes, increased acquisition times to counteract this SNR
457 loss.

458 Early accurate motion estimation techniques were developed using tagged MR and applied to
459 PET motion correction [147]. Tagged MR enforces a spatially periodic magnetization pattern (tags)
460 prior image acquisition; after the tagging pulses, the magnetization pattern is maintained over a
461 specific time and changed by motion, which can be measured in the MR images.

462 Several recent MR-based motion correction methods for cardiac PET data included specifically
463 designed MR protocols to account for motion during the examination [148–150]. Nevertheless,
464 these images present a lack of soft tissue contrast and/or spatial resolution that avoids using these
465 images for other tasks rather than motion estimation (i.e., for diagnostic purposes). Thus, the
466 inclusion of these additional images usually results in prolonged examination times, as MR images
467 acquired during the PET/MR study are just acquired for the sake of motion correction, and diagnostic
468 MR images have to be acquired once the PET acquisition is over.

469 Most current approaches have tried to avoid prolonging the MR acquisition protocol. On
470 the one hand, several methods have attempted to either extract motion estimates from respiratory
471 navigated 3D MR acquired images based on radial phase encoding [151], or from image-navigators
472 (iNAVs) to track cardiac movement being feasible in single phase whole heart imaging [152]. This
473 type of MR acquisition protocols demonstrated allowing for both high resolution 3D structural data
474 and extremely precise nonrigid motion estimation without an increase in overall scan times, leading
475 to a strong enhancement in both MR and PET image quality and guaranteeing a precise assessment
476 of radiotracer uptake. On the other hand, newer methods have tried to benefit from the use of CNNs
477 to estimate cardiac motion; for instance, CarMEN, an unsupervised 3D cardiac motion estimation
478 network for deformable motion estimation from 2D cine MR images [153]. This is an example on
479 how CNNs allowed for a novel framework where motion can be estimated from cine MR volumes,
480 despite each time frame being strictly a 2D stack of independently acquired sections, rather than an
481 actual 3D simultaneously acquired volume.

482 **4.4 Image reconstruction**

483 The goal of image reconstruction in PET is to obtain the most accurate representation of the
484 radiotracer distribution inside the patient. Different methods have been proposed over the years in
485 order to translate the emission-detected events into the 3D image space. Most authors have separated
486 reconstruction methods into analytical and iterative (or statistic) methods [154–156], although it

487 is worth mentioning the recent incorporation of DL techniques either for direct reconstruction or
488 mixed reconstruction with existing methods [157–160].

489 **4.4.1 Analytical methods**

490 A fundamental tool of these methods is the Radon transform [161], which defines the projection
491 space (also called sinogram in PET), where the PET data are found. To recover the image from the
492 projection space, the central-section theorem (or Fourier slice theorem) plays a major role. This
493 theorem correlates the Fourier transform (FT) of each 1D projection with the radial profile through
494 the 2D FT of the image at the same angle [162]. The most widely used analytical method is the
495 filtered-backprojection (FBP) algorithm [155, 161, 162]. Other analytical methods, such as the
496 backprojection-filtering (BPF) [163], or 3D adaptations of the FBP [164] have been also developed.
497 Recently, Lopez-Montes et al. have introduced a reconstruction method based on a more accurate
498 model of the system response matrix (SRM) using a pseudoinverse transformation that provides
499 higher resolution than the traditional FBP [165, 166].

500 In general, analytical methods assume idealized linear image projection models of the system
501 with additive gaussian noise and complete data that can be mathematically inverted or pseudo-
502 inverted to recover the activity distribution. These assumptions are not good approximations for
503 PET scanners, resulting in artifacts and quality loss in the reconstructed images. For instance,
504 the physical effects that degrade the image quality are not modeled in the FBP, thus limiting the
505 resolution that can be recovered from this method. In [167], the sinograms were deconvoluted
506 to restore the lost resolution because of intrinsic PET effects. The problem as it has been stated
507 so far is ill-posed, and small changes in the data (for example, because of Poisson noise, random
508 coincidences, scatter events, etc.) may lead to big changes in the image. Low-pass filters have
509 been commonly used to remove high-frequency noise. A few common examples are the ramp
510 filter, the Hamming filter or the Shepp-Logan filter [162]. PET scanners are also conformed of
511 discrete detector units with gaps in the projection domain from two contiguous blocks, which result
512 in discrete data with small gaps. This problem is partially solved using "inpainting" techniques to
513 fill in the empty spaces, or other sophisticated strategies [168]. Most of the analytical methods
514 are derived from 2D formulation, whereas PET data are gathered in 3D. Although some authors
515 adapted the problem to 3D [164], it is very common to use 2D rebinning techniques followed by
516 2D reconstruction [165, 166, 169–172]. Analytical methods produce noisier images than iterative
517 methods. In contrast, it is important to remark that analytical methods are usually faster, being a
518 linear mapping between the data and the reconstructed images, and they are typically preferred for
519 performance assessment of PET scanners [173, 174].

520 **4.4.2 Iterative methods**

521 Iterative methods focus on the optimization of an objective function, searching for the image that
522 better matches the data through small steps. These algorithms include a SRM that relates the
523 data domain with the image domain, the objective function to be optimized, a statistical model
524 of the noise in the data, and an iterative algorithm that estimates the image evolution toward the
525 convergence of the optimal solution [163, 175].

526 The SRM is of great importance to obtain the best image quality for our system, and it
527 must accurately model the most relevant physical processes involved in the emission, transport,

528 and detection of radiation in our acquisition system [156]. The most important considerations
529 are the positron range, the non-collinearity of annihilation photons, photon interactions with the
530 patient body (attenuation and scatter), random coincidences, the scanner geometry, and the detector
531 response (such as the depth of interaction in the detectors (DOI), inter-crystal scatter, or intrinsic
532 efficiency). Usually, the matrix is factorized and attenuation correction, scatter events, random
533 coincidences, and the normalization (sensitivity of a given LOR) are corrected separately [176, 177].

534 SRMs have been estimated using either analytical, experimental, MC, or hybrid approaches.
535 Analytical approaches are based on geometrical projection models [54, 178–181], analytical models
536 of the physical processes [182–185], or convolutional kernels usually referred to as Point Spread
537 Function (PSF) [186]. The PSF is a widely used method for SRM modeling [187], consisting of
538 either image or projection blurring kernels applied as a convolution operation. These methods
539 are fast, and they allow for on-the-fly calculations that drastically reduce the storage requirements
540 of a precomputed SRM [156, 188, 189]. For this reason, they have been usually combined with
541 experimental or MC estimations of the SRM to optimize the computational performance. In
542 experimental approaches, point sources are acquired at different positions of the scanner FOV and
543 the measured response is fitted to hand-crafted convolutional kernels either in the image [190–192]
544 or sinogram domain [193–197]. In [198], the authors used narrower PSF for single crystal events
545 identified by the detectors to increase image resolution. Other authors have benefited from MC
546 methods to reproduce the SRM with long simulations. In this case, the SRM may be determined as
547 an explicit matrix of probabilities relating to each voxel-LOR pair [199–203]. Since these matrices
548 require high loads of memory storage, symmetries and SRM sparsity have been exploited to reduce
549 the SRM size. Some authors applied PSF descriptions [185, 204–207], or wide tubes-of-response
550 (TORs) [208–211] instead of simple projection models. Other approaches include separate positron
551 range correction, which is included in the reconstruction as an image blurring kernel [212, 213].

552 The objective function and the optimization method used also play an important role in the final
553 solution achieved. The first proposed algorithms, known as Algebraic Reconstruction Techniques
554 (ART) [214], minimize the least-squares problem of the system. Later, a variation of the ART
555 algorithm was proposed by Andersen [215], the Simultaneous Algebraic Reconstruction Technique
556 (SART). One of the most successful algorithms has been the Maximum Likelihood Expectation
557 Maximization (MLEM) [216]. Motivated by the stochastic nature of the emissions in PET and
558 Single-Photon Emission Computed Tomography (SPECT), the objective function, in this case, is
559 the log-likelihood of the Poisson model for the data and the projections, and the optimizer chosen is
560 the expectation maximization [217]. The Image Space Reconstruction Algorithm (ISRA) [218] was
561 presented to preserve non-negativity images and reduce the memory requirement of the MLEM. The
562 MLEM was optimized with the Ordered Subsets EM (OSEM) algorithm [219], which reduces the
563 reconstruction time by the use of a small set of data in every update. OSEM is the current standard in
564 PET imaging. The use of subsets can be generalized to other optimizers to guarantee convergence,
565 like [220] did for the ascent gradient algorithm in the Row-Action Maximum Likelihood Algorithm
566 (RAMLA). The MLEM is presented as a particular case of RAMLA. In [221], both ISRA and
567 MLEM were presented from the point of view of a gradient descent algorithm of a weighted
568 least-squares objective function, with properly chosen descent steps. In terms of noise models,
569 the weighted least squares can be viewed as the log-likelihood of a gaussian model of the noise.
570 Inspired by the use of gradient-descent Stochastic Variant Reduction (SVR) algorithms in machine

571 learning [222], many authors have recently studied the application of these methods in PET imaging
572 to avoid limit cycles when ordered subsets are implemented [223–227].

573 In general, iterative methods converge to noisy images if the number of iterations is not
574 controlled. To avoid this, penalty functions (also called priors) are added to the objective function.
575 The goal is to select the image that satisfies our expectations, preserving edges while imposing local
576 smoothness. These methods are usually known as Bayesian or Maximum "a posteriori" (MAP)
577 methods [228, 229]. One of these implementations is the Block Sequential Regularized Expectation
578 Maximization (BSREM) [230, 231], presented as a globally convergent algorithm incorporating
579 priors in RAMLA with the appropriate step size. Examples of the penalty functions are the quadratic
580 function or the Huber functions [232], the Geman function [233], the relative differences prior (RDP)
581 [233–235], the median root prior (MRP) [236, 237], total variation [238, 239], or patch-based priors
582 [240]. In general, these methods are susceptible to proper hyperparameter selection, and their use in
583 the clinic is still limited. Q.clear [241–243] is one of the currently available commercial software,
584 based on BSREM and the RDP prior. In the Kernelized Expectation Maximization (KEM) method
585 [244], the image is modeled with a linear combination of specific kernels based on features that can
586 be obtained from static images in dynamic frames [244–246] or anatomical images [247, 248].
587 KEM allows a more complex representation of the priors since spatially-variant dependence can be
588 included, instead of fixed functions for the whole image. This method has also been proposed as a
589 successful choice to incorporate DL within the image reconstruction process (see next section).

590 All of these components make iterative algorithms outperform analytical algorithms in terms
591 of SNR [154]. The statistical model makes them more robust to noise and incomplete data,
592 and together with the iterative scheme allows for non-linear considerations in the reconstruction
593 workflow. They can also introduce geometrical properties and model physical processes in the SRM,
594 thus making them more accurate and easily adaptable for 3D imaging [155]. Iterative methods are
595 computationally intensive compared with the FBP, since each iteration requires a forward projection
596 and a backward projection operation, whereas in FBP this step is performed only once. Nevertheless,
597 nowadays computer power and GPU parallel computing allow for routine image reconstruction with
598 iterative methods in short computation time [203, 249–251].

599 4.4.3 Deep learning methods

600 In image reconstruction of PET data, the application of DL has followed different strategies [157,
601 158]. So far, we can summarize them in direct reconstruction, physics-informed DL, and post-
602 processing methods [160, 252].

603 Direct reconstruction consists in mapping from noisy data to the output image, thus learning
604 the full transformation from the projection space to the image space [253–257]. This approach
605 currently implies memory issues since both spaces are high dimensional, requiring reducing the
606 problem to individual 2D slices and large data sets to train the model.

607 In physics-informed DL methods, a hybrid approach combining DL with known physics about
608 the system (such as the projection/backprojection model, the PSF, etc.) is implemented. Different
609 authors implemented a fixed backprojection step combined with learned filters [258] and an image
610 postprocessing network [259]. In [260–262], a hybrid approach was implemented for TOF data
611 using histograms based on the most-likely annihilation position, and [263] implemented a cycle-
612 consistent learning-based hybrid iterative reconstruction method based on two networks mapping

613 from the backprojection to the reconstructed image and vice versa. These approaches may be
614 classified as direct DL reconstruction methods by some authors since they minimally make use
615 of the Radon transformation or simple backprojection instead of the traditional iterative methods.
616 Other authors have embedded DL methods within image regularization in the image reconstruction
617 process. One of the methods proposed is the implementation of a learned gradient prior in the
618 MAP algorithm [264–268]. The Deep Image Prior (DIM) method introduced by Ulyanov et al.
619 [269], which consisted in the use of a CNN without specific training for image denoising, was
620 incorporated in the ML optimization by Gong et al. using MR anatomical information [270]. DIM
621 has the advantage of using unsupervised learning based on the patient information to train the
622 network during reconstruction, thus not requiring large training datasets. Other authors have also
623 studied the implementation of DIM within the reconstruction [271, 272]. A CNN has also been
624 incorporated within the KEM method [273], and eventually, it has been combined with DIM [274].

625 The third approach is performing image post-processing with DL methods after reconstruction
626 to improve quality or denoising. One of the most attractive applications of DL is image denoising
627 from low-dose PET to standard-dose PET [275–280], as it may reduce the radiation exposure of
628 patients in PET studies. The advantage of this approach is that image pairs can be easily obtained by
629 removing events from list-mode data of standard-dose acquisitions. Denoising is also challenging for
630 gated PET images [281], where gates suffer from higher noise than static images and different noise
631 levels among gates. Image quality improvement has also been achieved including MR inputs in the
632 CNN training [282–285]. These approaches demonstrated that MR images helped to significantly
633 decrease noise while improving resolution and preserving detailed structures. Other authors applied
634 a CNN to correct positron range blurring in ^{68}Ga images [286, 287], or transforming non-TOF to
635 TOF images [288]. In the previous section 4.2, we already mentioned similar post-processing DL
636 approaches to correct attenuation and scatter in the image domain.

637 4.4.4 MR-guided and joint reconstruction

638 Several recent approaches have focused on using both classical image processing techniques as
639 well as novel techniques such as CNNs to improve PET or MR imaging. Therefore, numerous
640 methods appraised using the potential of simultaneous PET/MR imaging in procedures to improve
641 the quality and information content of original data while reducing artifacts, accelerate imaging
642 acquisition and reconstruction, or decrease the associated risk to the procedures.

643 Many proposals directed efforts on improving the quality and information content of original
644 data while reducing artifacts, especially as PET imaging suffers from severe inherent resolution
645 limitations which reduce its quantitative spatial accuracy and degrades PET images by partial
646 volume effects. In this context, one of the most studied approaches, both for PET and MR, is
647 super-resolution. Super-resolution imaging refers to a category of techniques that enhance the
648 quality of the acquired images by boosting their apparent resolution. In this scenario, CNN-based
649 super-resolution methods can generate high-quality images from those reduced-resolution ones. In
650 this sense, we can find methods for intramodality image super-resolution in PET imaging [289],
651 in MR imaging [290], but increasingly more taking advantage of the complementary information
652 present in combined PET/MR images [291]. The latter methods proved how CNNs trained using
653 combined data outperformed traditional methods by large margins. Related to super-resolution
654 methods, we find denoising methods, which were already discussed in section 4.4.3.

655 Other methodologies have also centered on reducing MR acquisition and reconstruction times.
656 This is mainly because, in a combined PET/MR scan, the MR image acquisition can be very time-
657 consuming and take more time than the actual PET scan. Minimizing the time required for raw
658 data acquisition or under-sampling the k-space can accelerate the acquisition, but this impacts the
659 resulting quality of the MR images. In this sense, there still exist novel developments to improve
660 the quality of the resulting MR images based on probabilistic models and other information theory-
661 based methods [292, 293], while CNNs are also gaining leadership in dealing with these applications
662 [294]. These methods were able to enhance image resolution without any hardware adjustments
663 while decreasing overall scanning time.

664 Furthermore, in previous sections we have already discussed several other approaches focused
665 on applying DL methods to correct attenuation and scatter (section 4.1 and 4.2), and reduce
666 radiotracer requirements without sacrificing diagnostic quality (section 4.4.3) in integrated PET/MR
667 systems.

668 **5 Conclusions**

669 Hybrid medical imaging modalities combine two modalities in a single scanner, providing rich
670 anatomical and functional information in a single study, thus reducing burden to the patient, time
671 and cost. Among these devices, PET/MR is the most promising one for clinical routine, but
672 the combination of these two modalities has required complex technical developments, both in the
673 hardware and in the associated data and image processing techniques. In this work we have reviewed
674 the history of PET/MR development, describing the different technical issues that were confronted
675 and how they were solved, with a special emphasis on instrumentation, and image reconstruction
676 and correction.

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